

Идеи и методы физики конденсированного состояния

XIV Школа-конференция молодых ученых
"Проблемы физики твердого тела и высоких давлений"

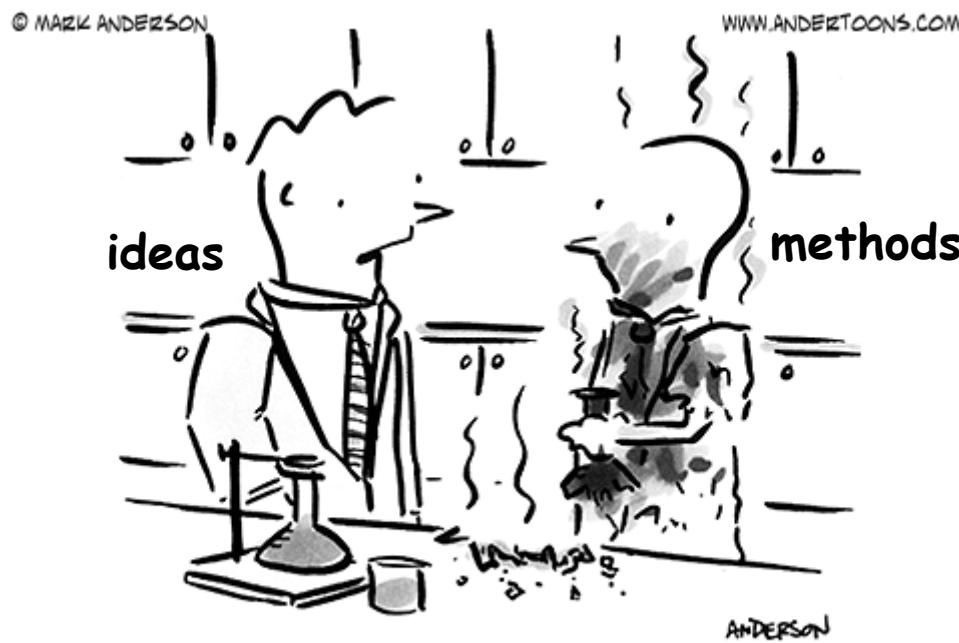
Сочи, 11-20 сентября 2015г.



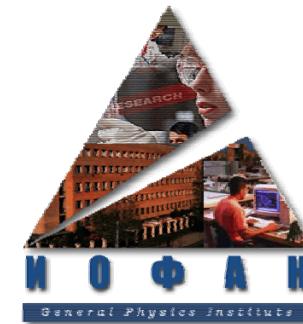
Magnetic resonance and quantum criticality

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"And therein lies the problem."



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Идеи и методы...

Неполиткорректная лекция

...Уступая моде, распространенной среди физиков первой половины XX века, интерпретировать все формально получаемые результаты с помощью простых физических представлений, выделим...

... Можно, конечно в духе классиков начала XX века еще порассуждать, например представить, что...

*Иридий Александрович Квасников,
II, 270*



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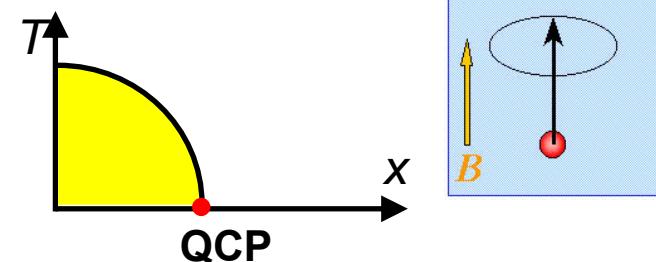
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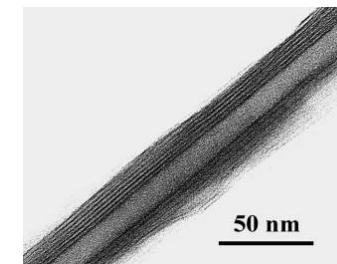
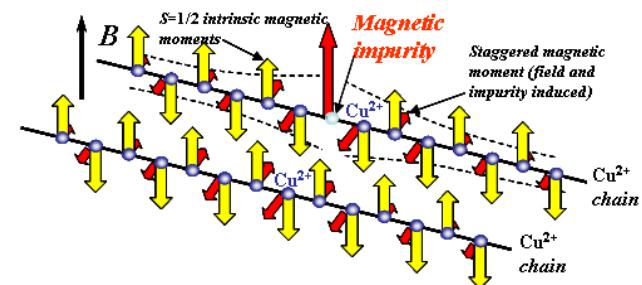
"Pavlov had his dog. Schrödinger had his cat. But nobody's done anything with a *goldfish* yet!"

If there any regularities which link quantum critical phenomena and magnetic resonance?

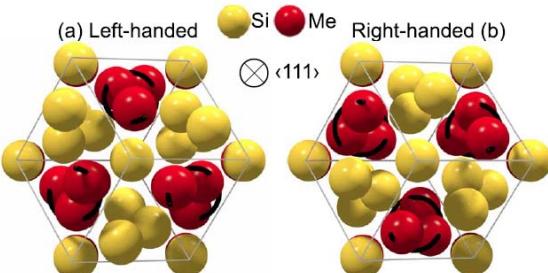
Introduction. Quantum criticality and electron spin resonance



Quantum spin chains in disorder driven quantum critical regime.
(Dielectrics, 1D systems).



Quantum critical phenomena in the nano-world.
(Bad conductors, 2D systems).



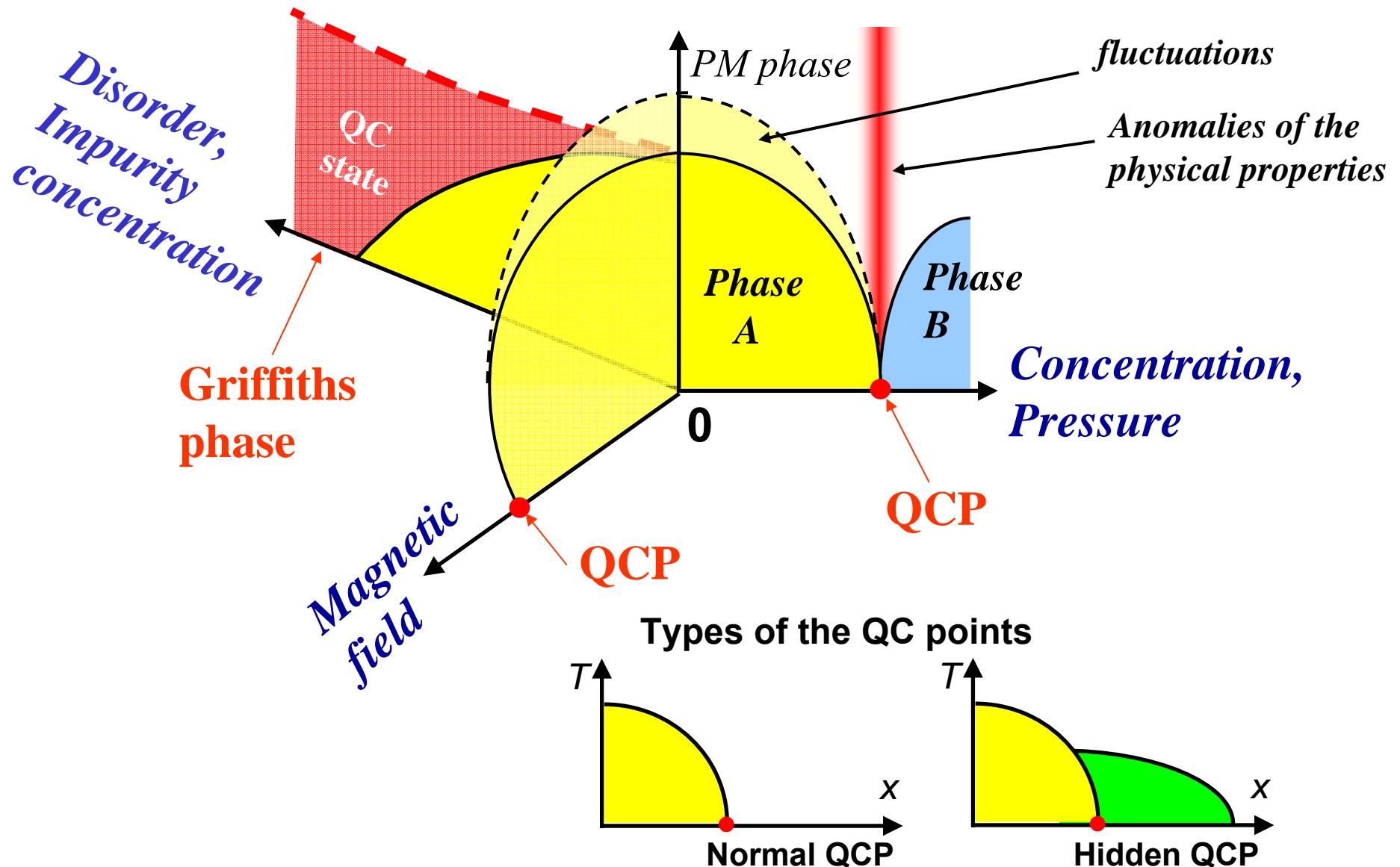
Quantum criticality in strongly correlated metals.
(Good conductors, 3D systems)

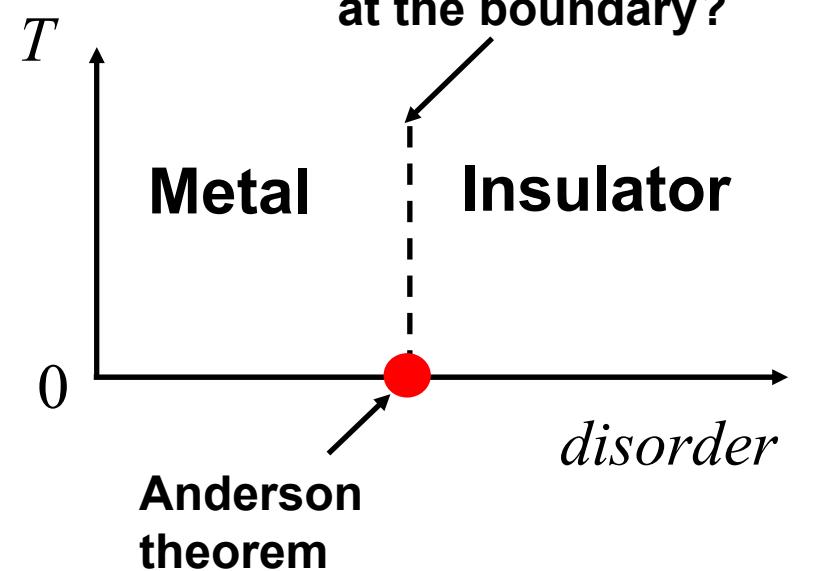
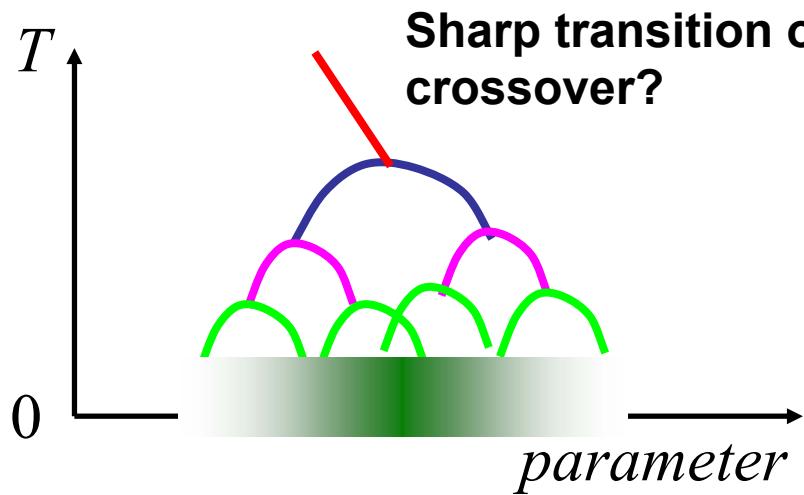
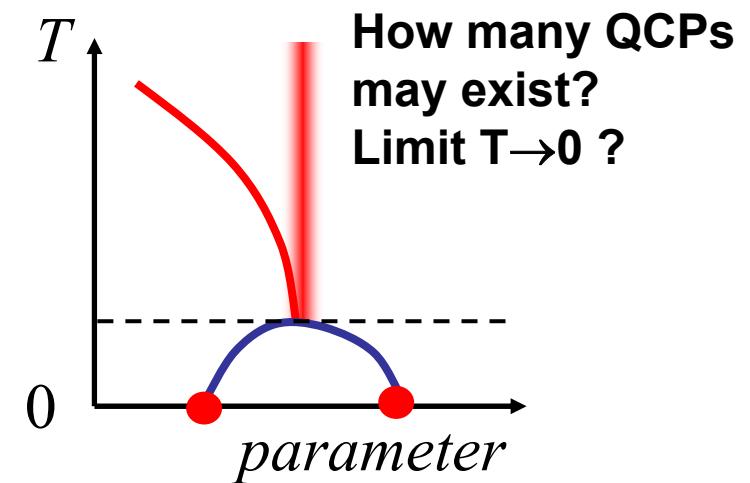
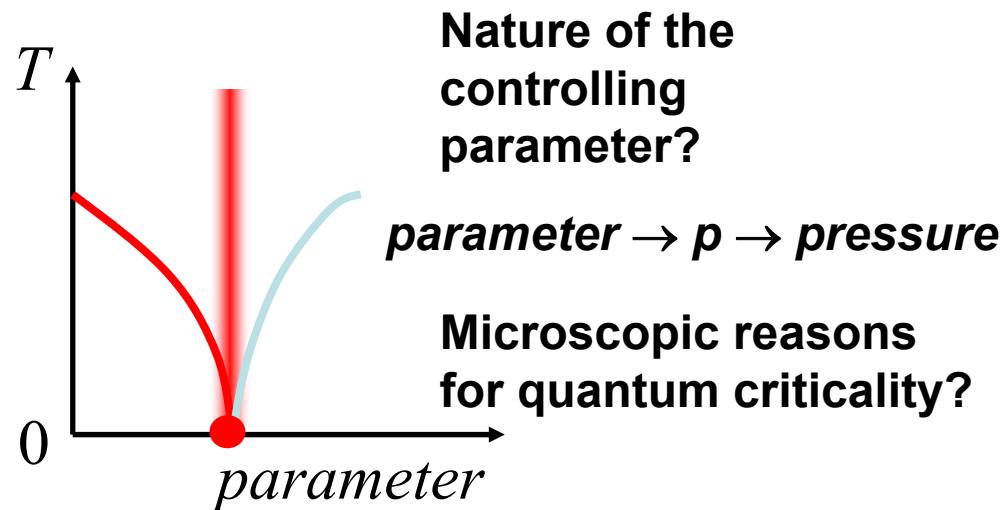
Final remarks



Quantum phase transitions are phase transitions at $T=0$.

Temperature

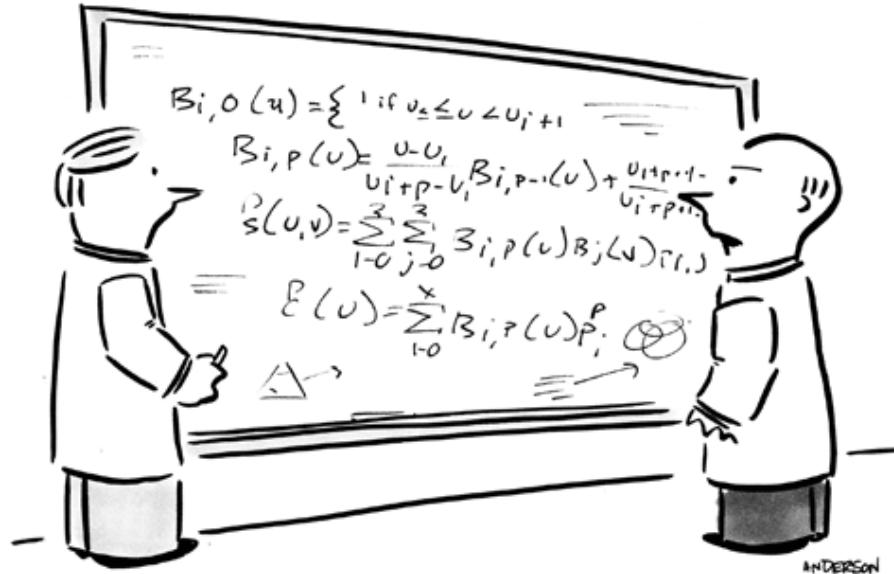






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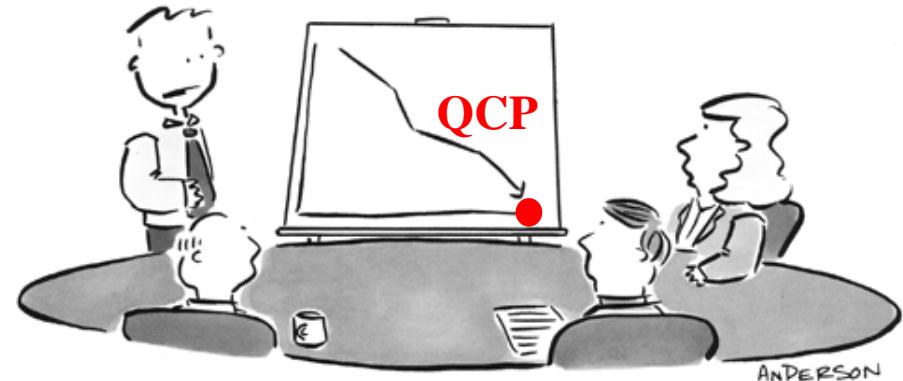
"What the hell is *that* supposed to mean?!"

Theory...

Experiment...

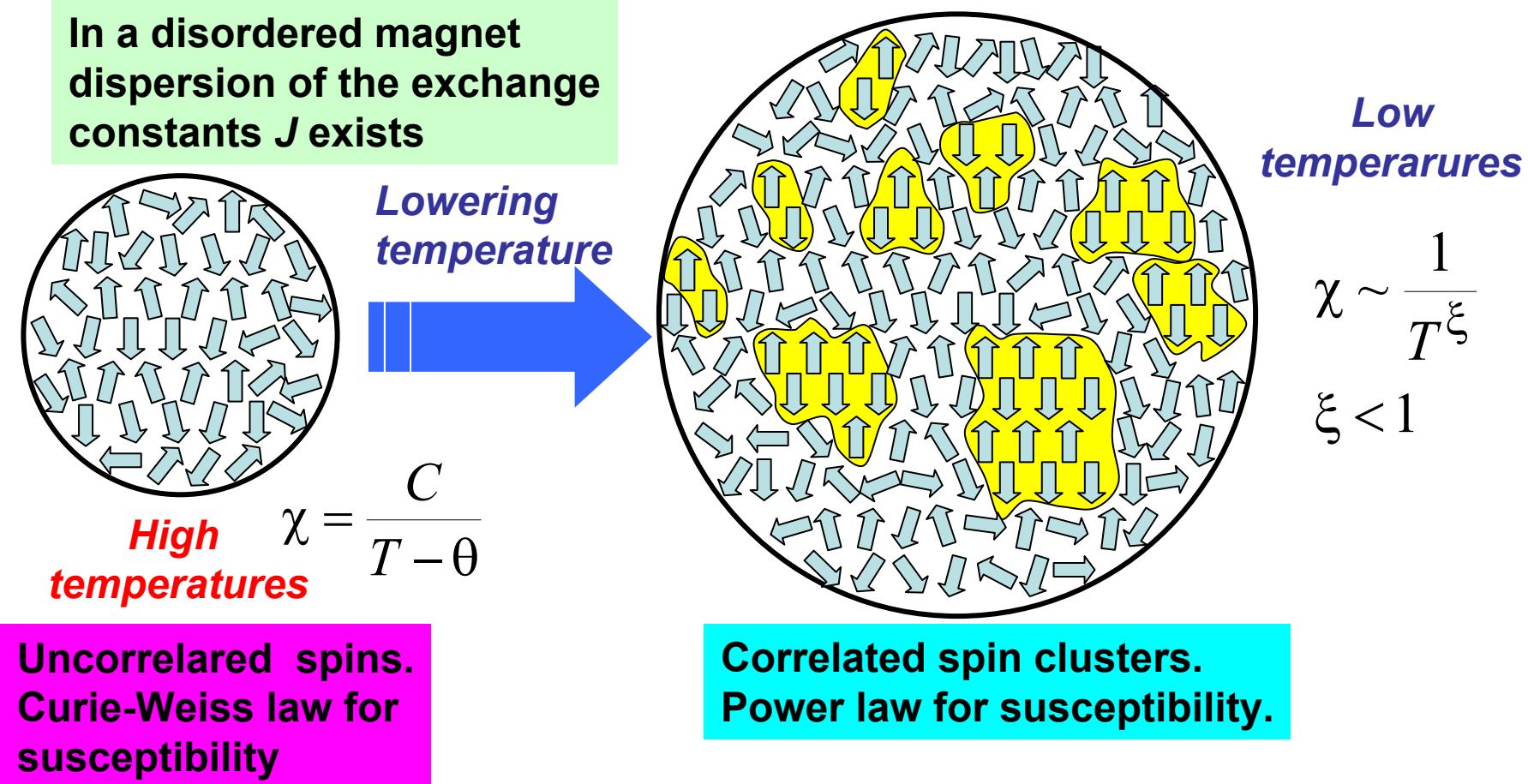
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"As you can see, misery loves our company specifically."

"Please, be so kind to tell us what kind of transition do we have."

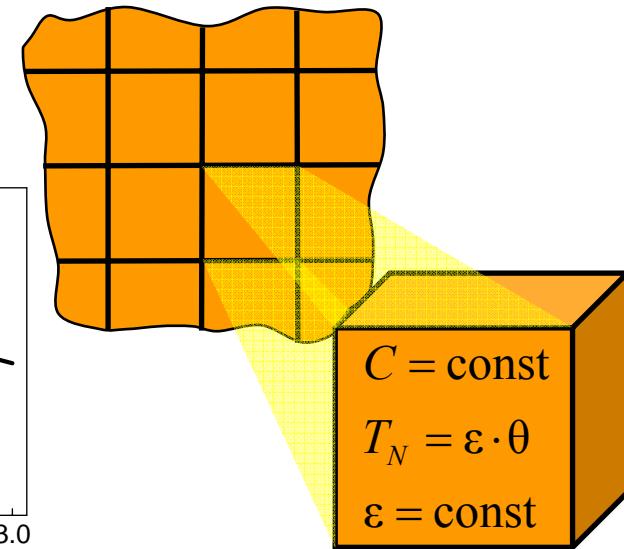
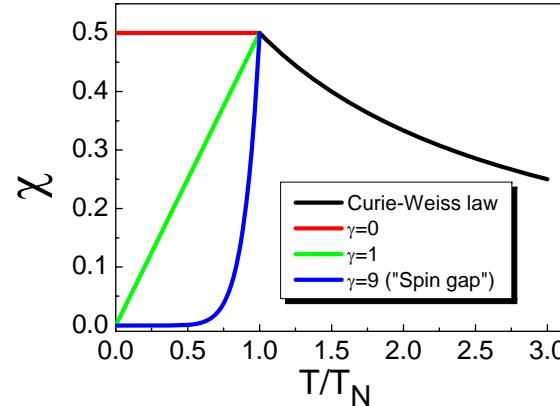
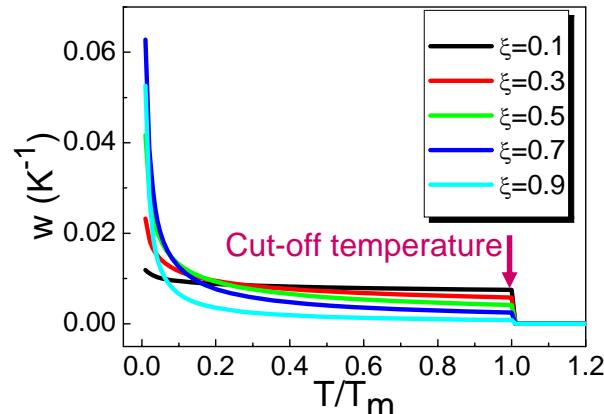


R.B.Griffiths. Phys. Rev. Lett., **23**, 17 (1969); A.J.Bray, Phys. Rev. Lett., **59**, 586 (1987)



The model:

S.V.Demishev, Phys.Sol.St., **51**(3), 514 (2009)
 S.V.Demishev, Phys. Stat. Sol. B, **247**(3), 676 (2010)



Distribution of Neel temperatures in the sample:

$$w(T_N) = (1 - \xi)(T_m / T_N)^{\xi} / T_m$$

Result (susceptibility at an arbitrary temperature):

$$\chi = C / (T + \theta^*) \quad (T > T_m)$$

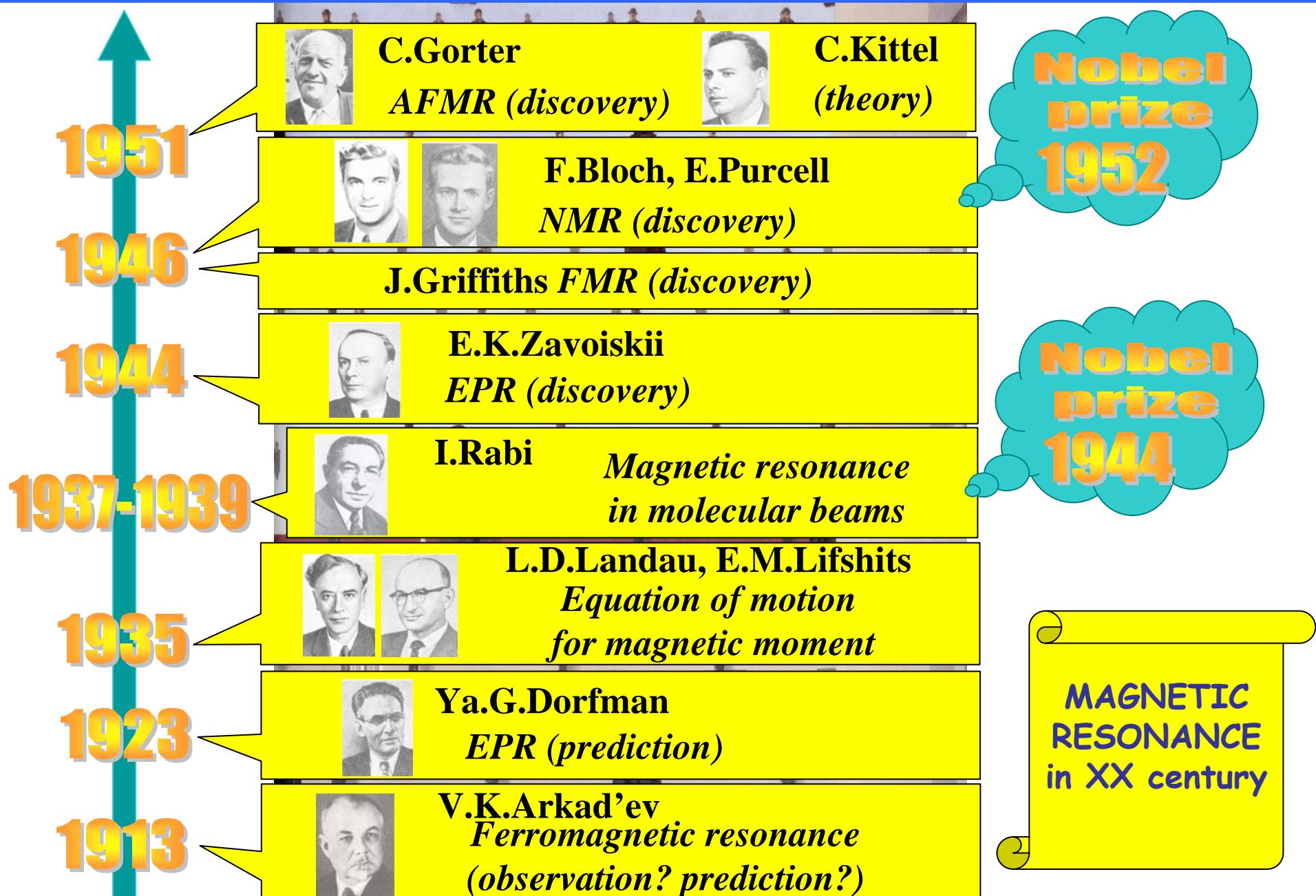
$$\chi = [C / (T_m + \theta^*)] (T / T_m)^{-\xi} \{1 + (1 - \xi)(1 + \theta^* / T_m)[1 - (T / T_m)^{\gamma + \xi}] / [(1 + \varepsilon^{-1})(\gamma + \xi)]\} \quad (T < T_m)$$

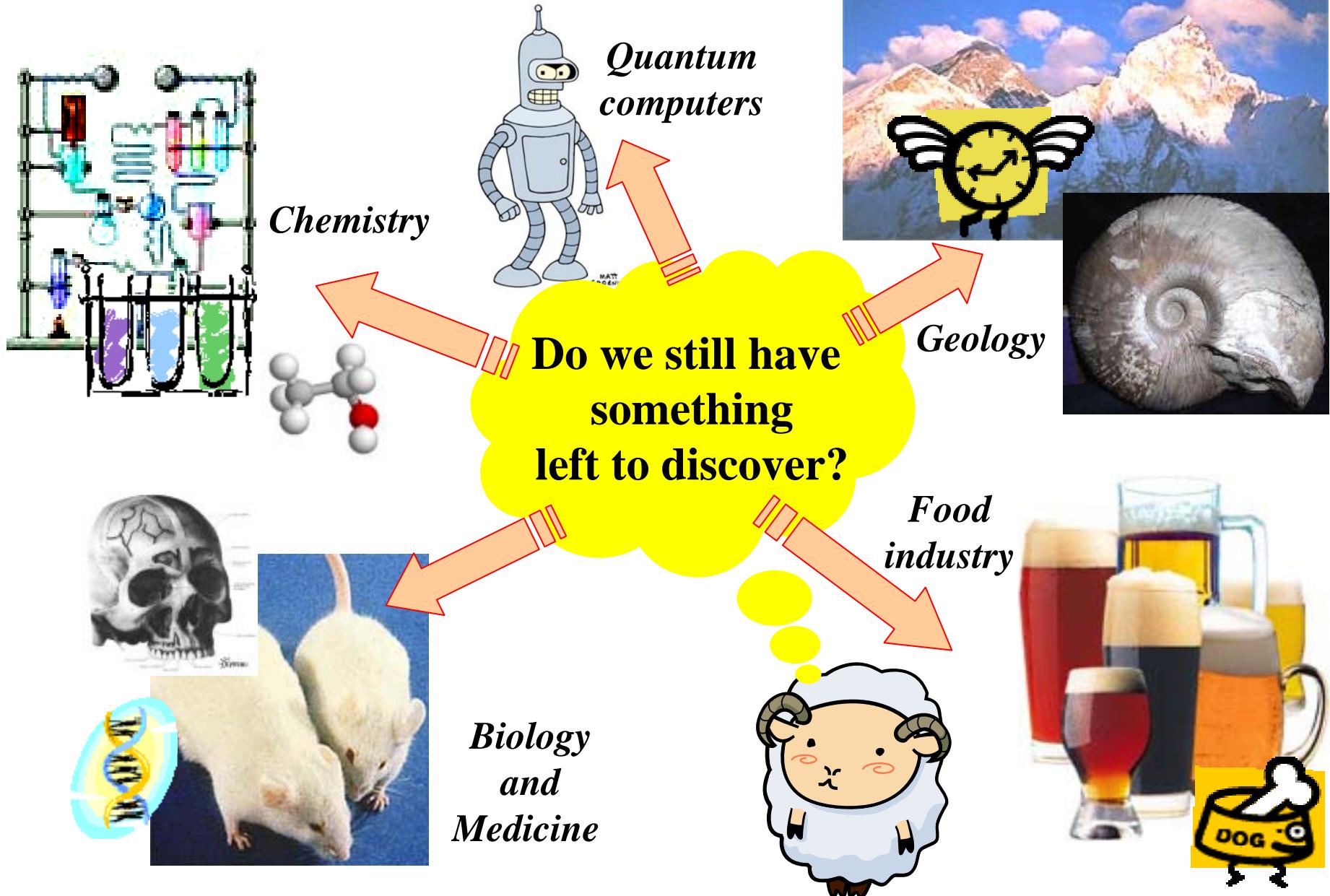
$$\theta^* = T_m [\varepsilon(1 - \xi)f(\varepsilon, \xi)]^{-1} - 1 \quad f(\varepsilon, \xi) = -\xi^{-1} - \sum_{r=1}^{\infty} (-1)^r \varepsilon^r / (r + \xi) + \pi \varepsilon^{-\xi} / \sin[\pi(1 - \xi)]$$

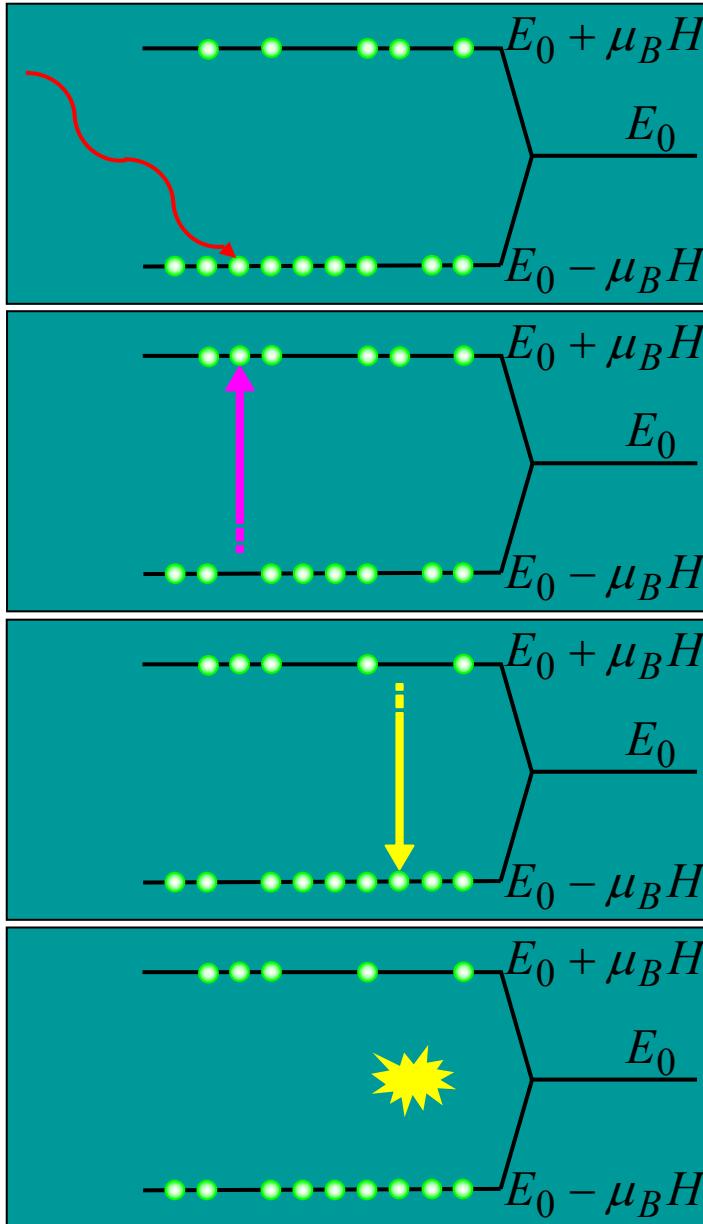
Susceptibility of the cluster with T_N :

$$\chi = C / (T + \varepsilon^{-1} T_N) \quad (T > T_N)$$

$$\chi = C (T / T_N)^{\gamma} / [T_N (1 + \varepsilon^{-1})] \quad (T < T_N)$$







Radiation absorption

Dynamic susceptibility tensor

$$\vec{H} = \vec{H}_0 + \vec{h}$$

$$\vec{M} = \vec{M}_0 + \vec{m}$$

$$\vec{m} = \hat{\chi}(\omega, B_0, M_0) \cdot \vec{h}$$

$$\vec{h}, \vec{m} \sim \exp(i\omega t)$$

Difficult to calculate in general case!

Quantum transitions

But... There is a simple
classical language...

Energy dissipation



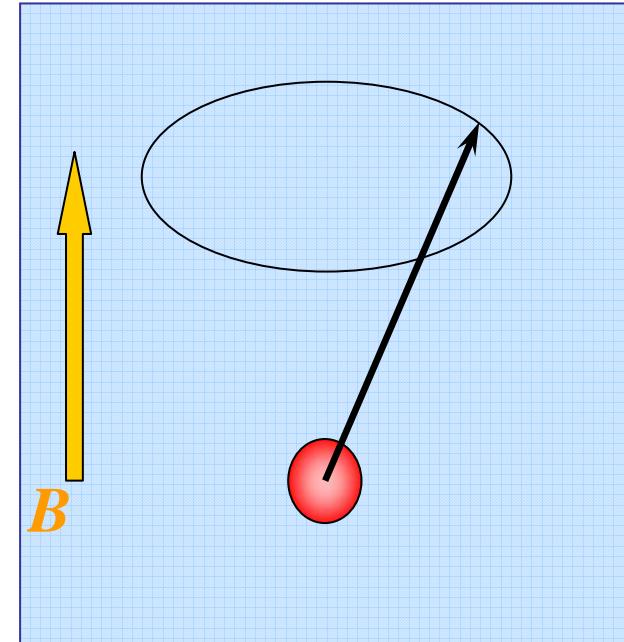


Spin = magnetic arrow

*Precession of a magnetic moment
in magnetic field*

Landau-Lifshits equation of motion:

$$\frac{d\vec{M}}{dt} = -\gamma \cdot \vec{M} \times \vec{H}_{loc} + \vec{R}(\vec{M})$$



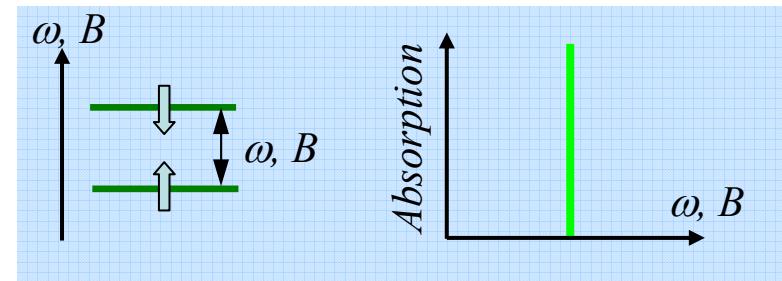
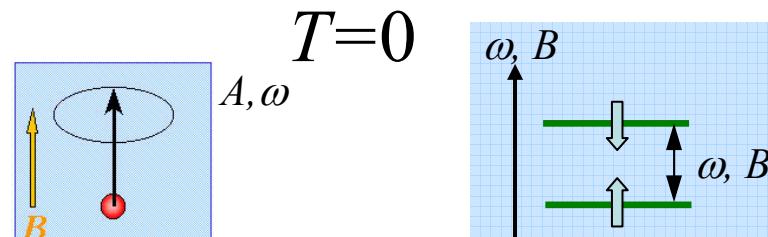
$$\hat{\chi} = \begin{pmatrix} \chi & i\chi_a & 0 \\ -i\chi_a & \chi & 0 \\ 0 & 0 & \chi_{II} \end{pmatrix}$$

$$\vec{R} = \frac{\alpha}{M} \cdot \vec{M} \times \frac{\partial \vec{M}}{\partial t}$$

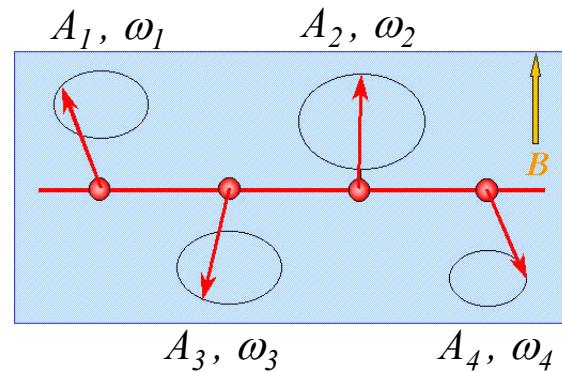




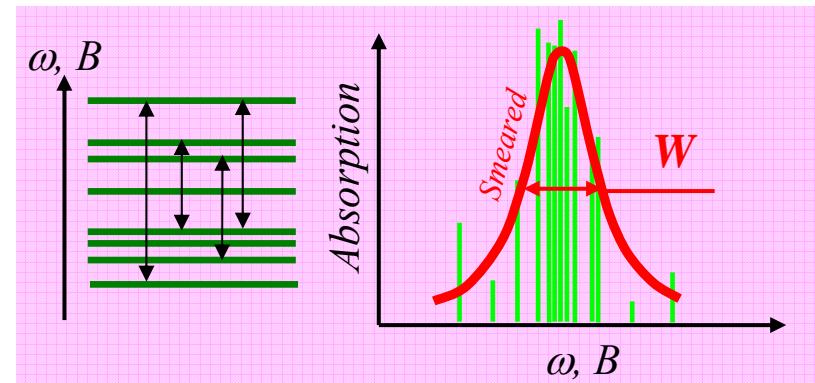
Single spin system



Strongly correlated spin system



$T=0$



$T \neq 0 ??$

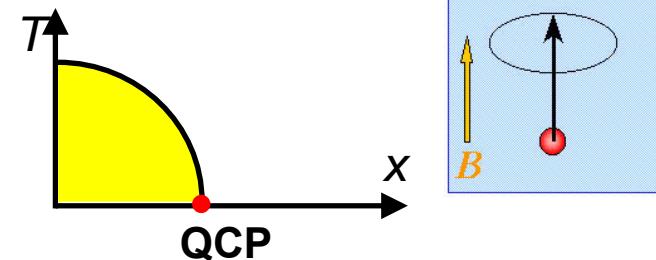
Quantum description is difficult and possible in a limited set of cases
(AF S=1/2 quantum spin chain, Oshikawa-Affleck theory).

Classical spin rotation in quantum critical phenomena?

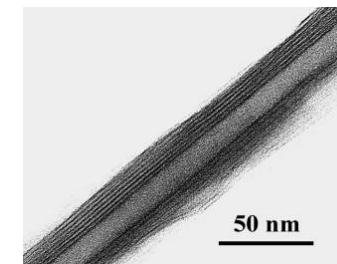
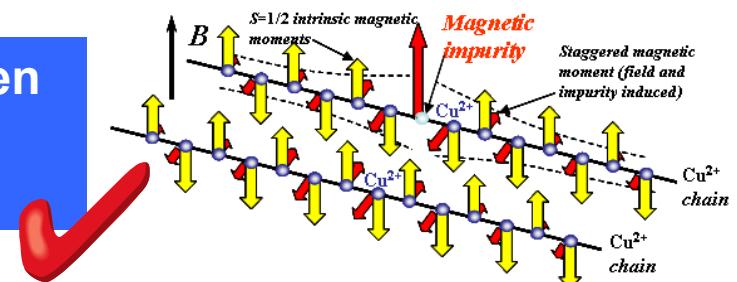


		We can	No way
Isolated spin	Fine structure of absorption line		
	Line shape		Sometimes...
	Selection rules		Incorrectly...
	Experimental geometry, polarization		If you are clever...
	Strongly correlated system	AF S=1/2 chain Abrahams-Wolffle theory	Easily!

Introduction: Quantum criticality and electron spin resonance

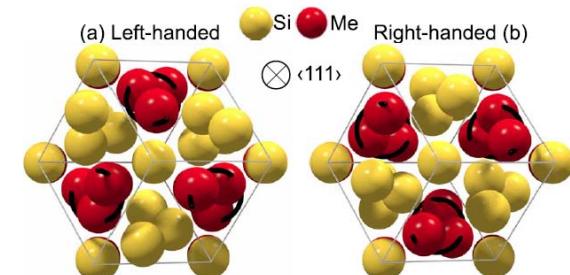


Quantum spin chains in disorder driven quantum critical regime.
(Dielectrics, 1D systems).

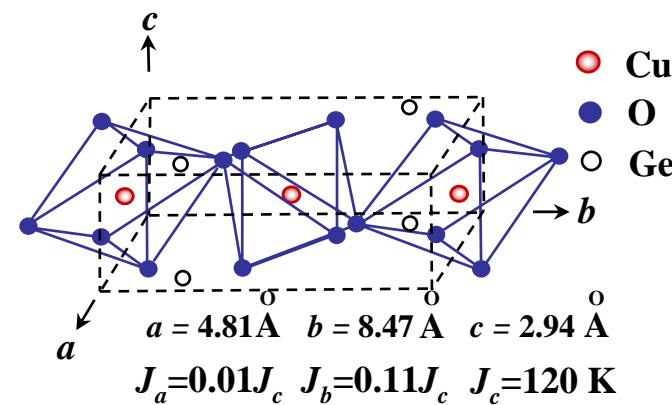
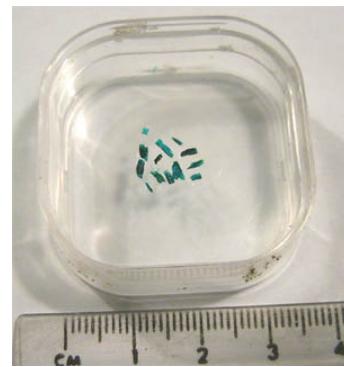


Quantum critical phenomena in the nano-world.
(Bad conductors, 2D systems)

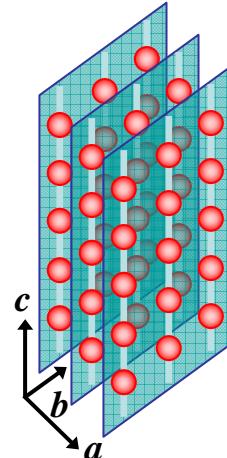
Quantum criticality in strongly correlated metals.
(Good conductors, 3D systems)



Final remarks



$S=3/2$ [Co] $S=2$ [Fe] $S=5/2$ [Mn]

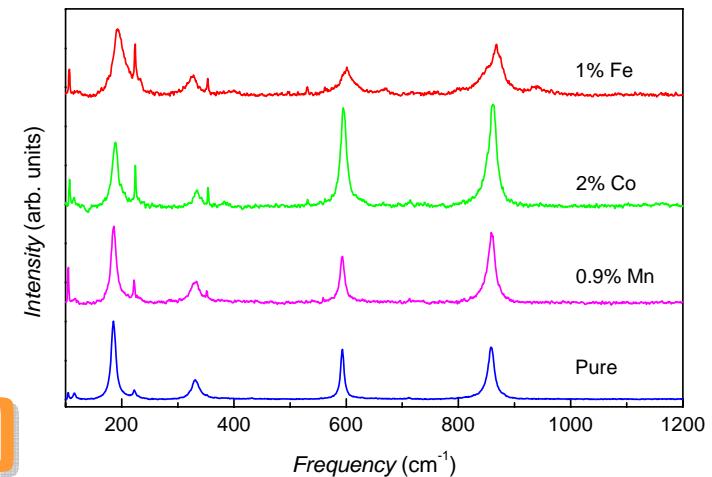


- Doping of $S=1/2$ (Cu^{2+}) chains with $S=3/2$ (Co^{2+}) $S=2$ (Fe^{2+}) and $S=5/2$ (Mn^{2+}) impurities
- Effect of doping-induced structural disorder

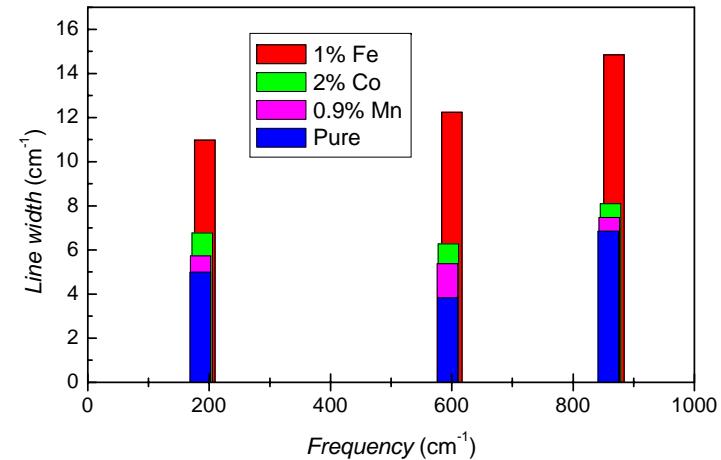
$S=1/2$ [Cu] AF chains

$\text{Cu}_{1-x}\text{M}_x\text{GeO}_3$ ($x=0.02$ M=Co; $x=0.01$ M=Fe, $x=0.009$ M=Mn)

Raman spectra



Line width and position



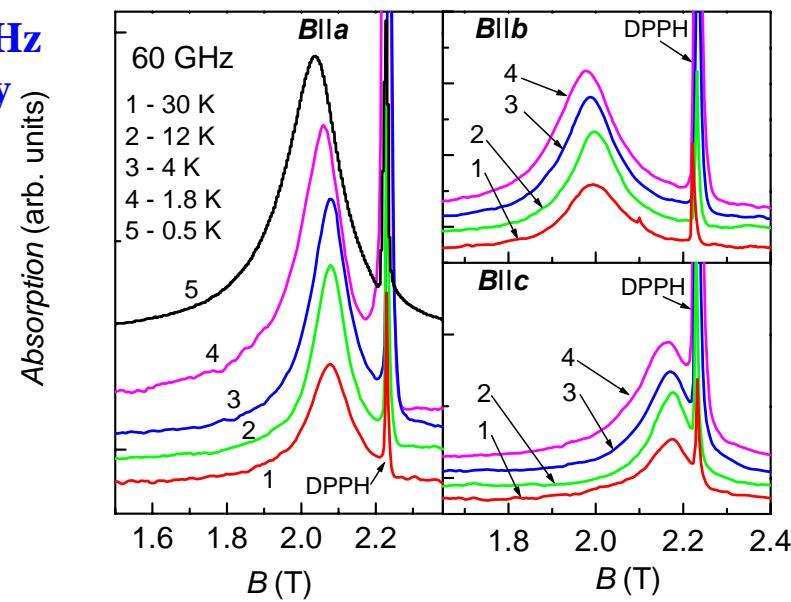


Experimental ESR spectra.

13

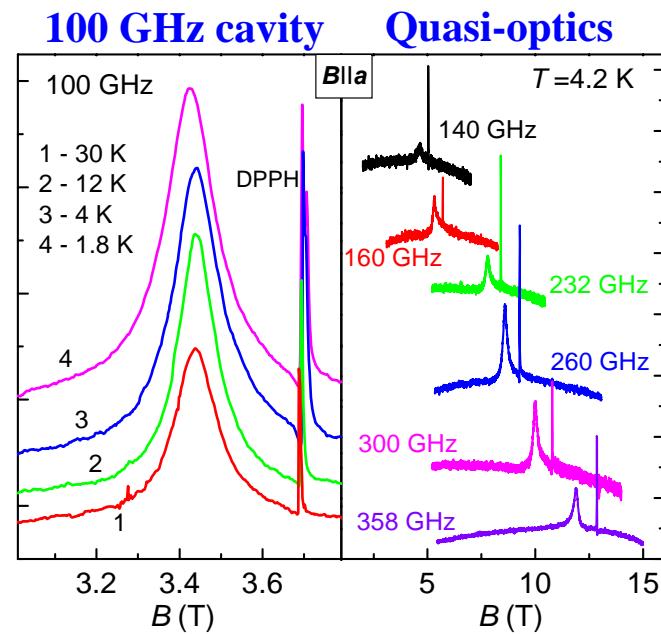
60 GHz
cavity

$S=2$ [Fe]



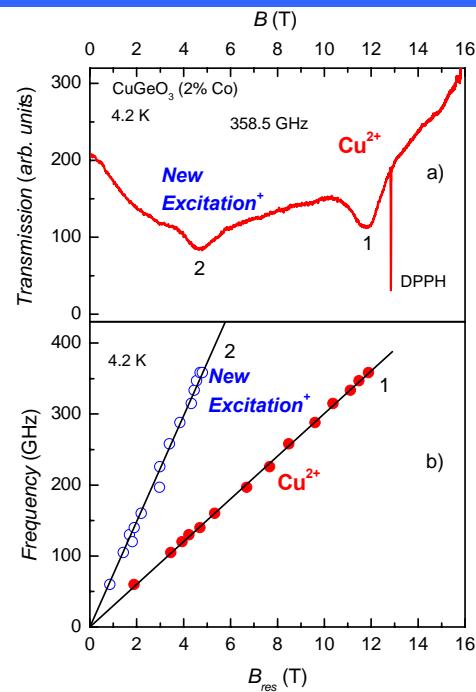
100 GHz cavity

Absorption (arb. units)

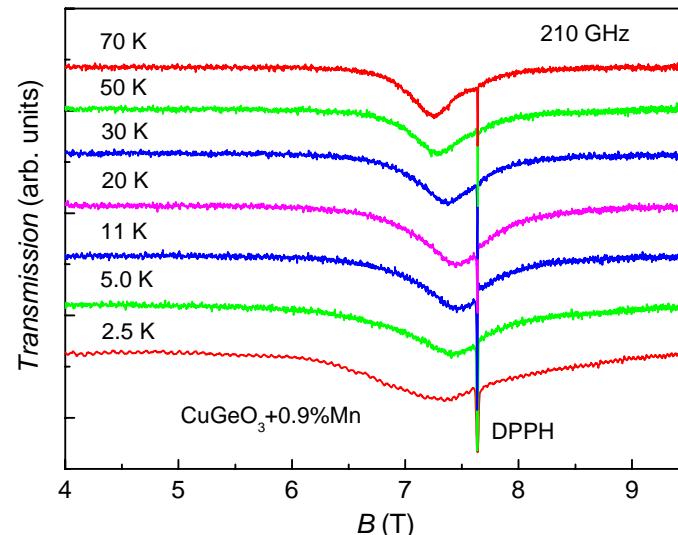


Quasi-optics

$S=3/2$ [Co]



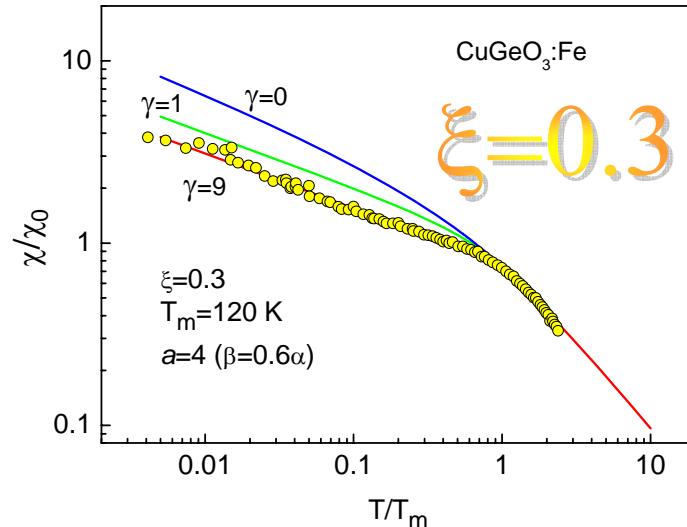
Quasi-optics



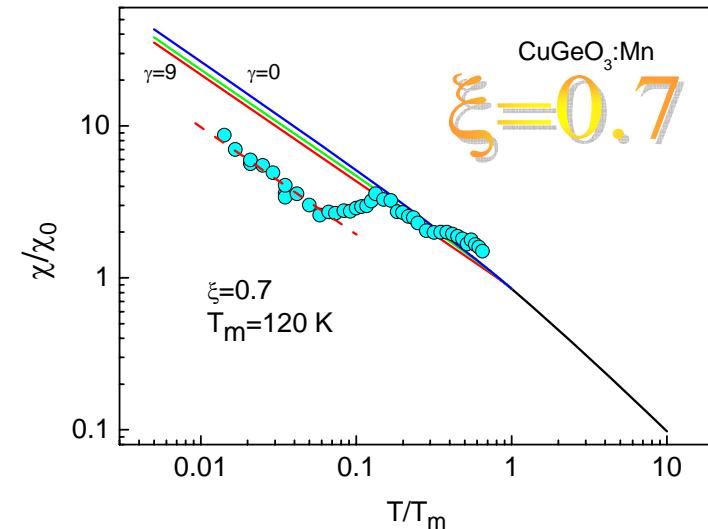
$S=5/2$ [Mn]



S=2 (Fe)



S=5/2 (Mn)



ESR monitors quantum spin chains

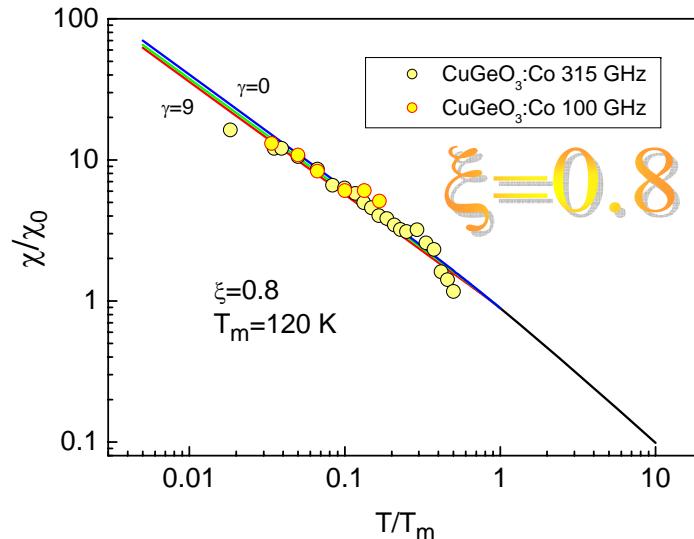
Spin-Peierls transition is suppressed by doping

In QC regime temperature vary more than two orders of magnitude.

Theory works!

$T_m \sim 120 \text{ K} \sim J_c$

S=3/2 (Co)





$$\hat{H} = \hat{H}_0 + \hat{H}_Z = J \cdot \sum_{j,k} \vec{S}_j \cdot \vec{S}_k - B \cdot \sum_j S_j^z$$

$$\hat{H} = \hat{H}_0 + \hat{H}_A + \hat{H}_Z$$

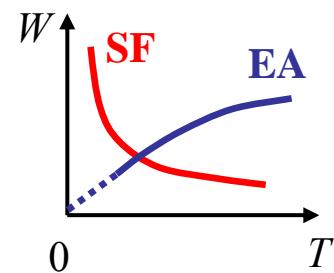
$$\hat{H}_A = h \cdot \sum_j (-1)^j \cdot S_j^x$$

Staggered field

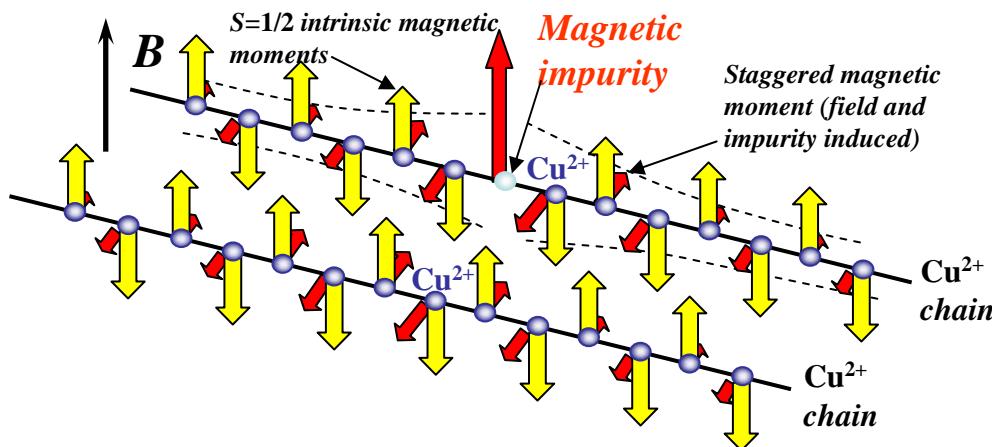
$$\hat{H}_A = \delta \cdot \sum_j S_j^x \cdot S_{j+1}^x$$

Exchange anisotropy

M.Oshikawa, I.Affleck (1999- 2002)



Line width in presence of exchange anisotropy and staggered field



Staggered magnetization is a fingerprint of QC regime in CuGeO₃?

Universal relation in OA theory

$$\frac{w}{\Delta g} = 1.99 \frac{k_B}{\mu_B} \cdot T$$

S.V.Demishev et al., *Europhysics Lett.* (2003)

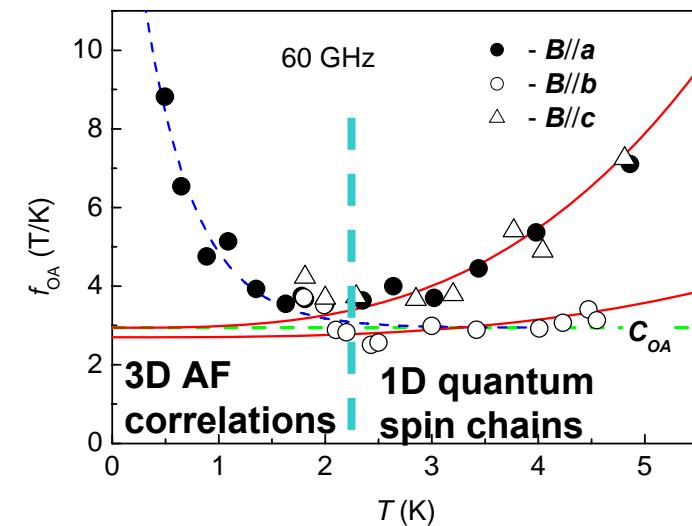
S.V.Demishev et al., *ProgTheor.Phys.Suppl.*(2005)

Oshikawa-Affleck function:

$$w(T) = a \cdot T + \gamma \frac{h^2}{T^2}; \quad \Delta g(T) = \zeta \cdot \frac{h^2}{T^3}$$

$$f_{OA}(T) \equiv \frac{w(T)}{\Delta g(T) \cdot T} \approx C_{OA} + \frac{a}{\zeta \cdot h^2} \cdot T^3$$

$$C_{OA} = 1.99 \cdot k_B / \mu_B$$

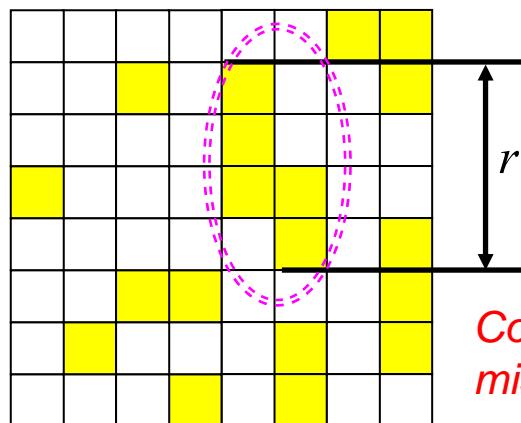




Lowering temperature results in freezing of the magnetic contribution of the spin clusters for which condition $T_N > T$ holds.

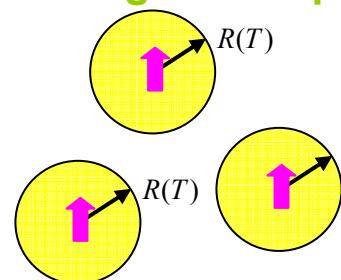
A two-phase model (“frozen” and “non-frozen clusters”) may be introduced. The probability $W(T_N)$ defines “frozen” part of the sample volume.

Percolation



Correlation is missing

Cluster growth in vicinity of the magnetic impurity



Correlation is essential

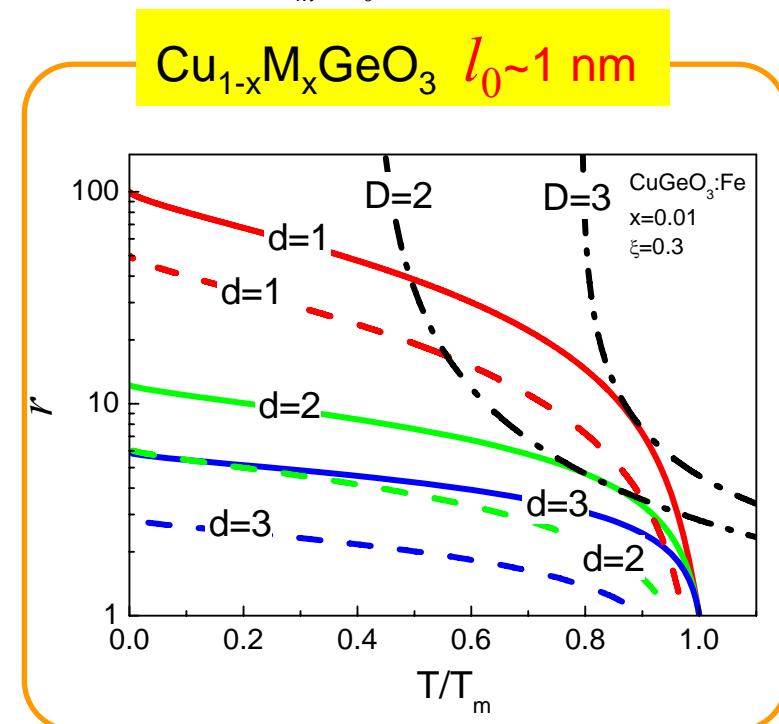
$$r(T) \sim \frac{1}{|v_{AF}(T) - v_c|^\eta}$$

$$\frac{T}{T_m} = \left[1 - \frac{2x}{d+1} \left(r^d - \frac{1}{r} \right) \right]^{\frac{1}{1-\xi}}$$

$$v_{AF} = 1 - \left(\frac{T}{T_m} \right)^{1-\xi}$$

$$R < R_m$$

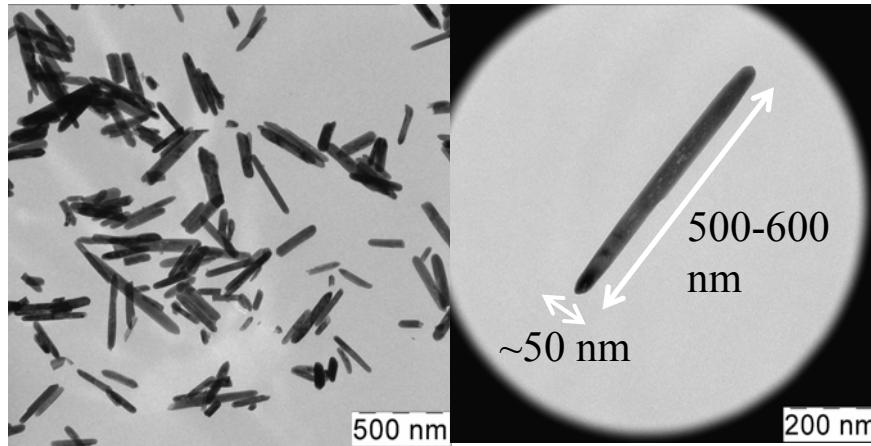
$$r = R_m / l_0$$



Griffiths phase is formed by nano spin clusters!

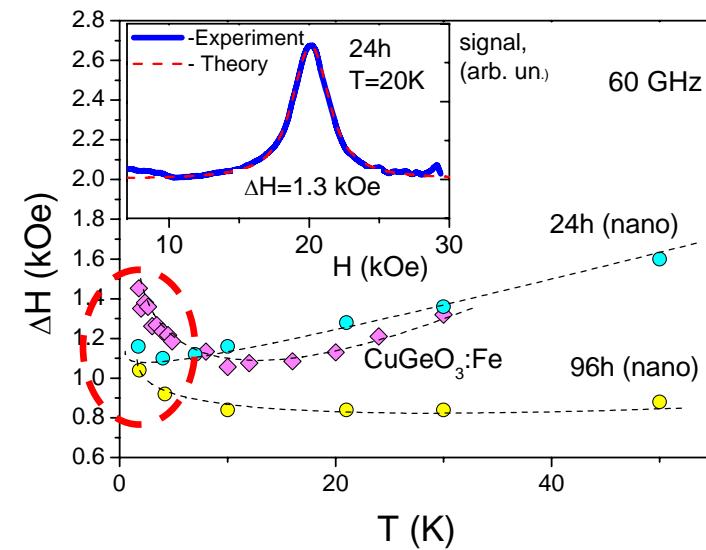
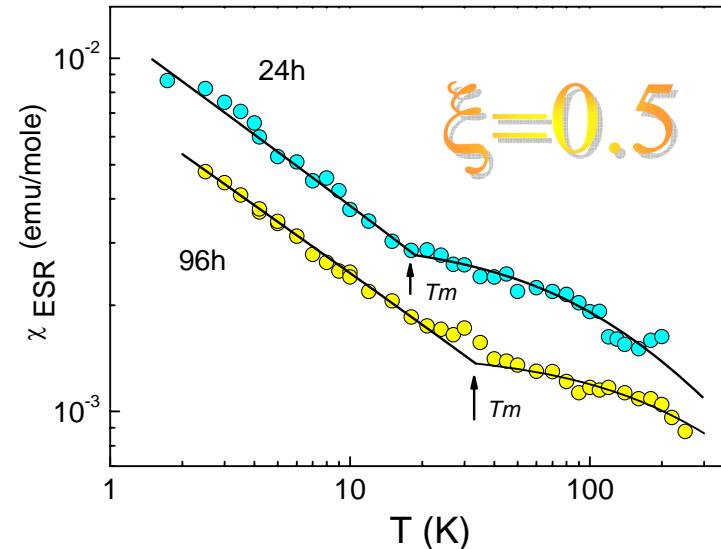


Synthesis by hydrothermal treatment

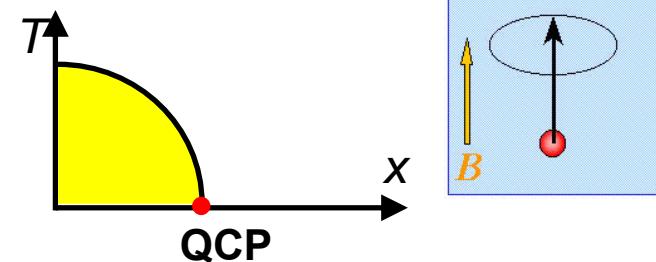


Suppression of spin-Peierls transition and QC behavior without any magnetic impurities!

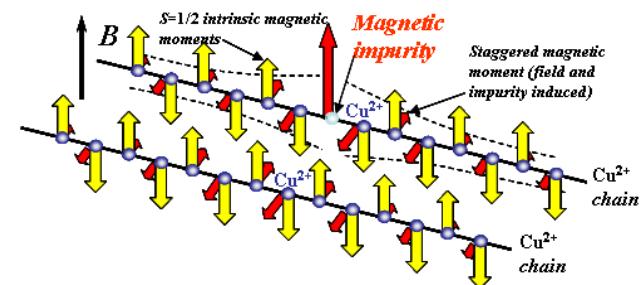
Line width simulation for a mixture with anisotropic g-factors suggests presence of the staggered field.



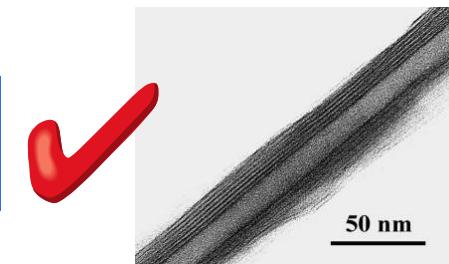
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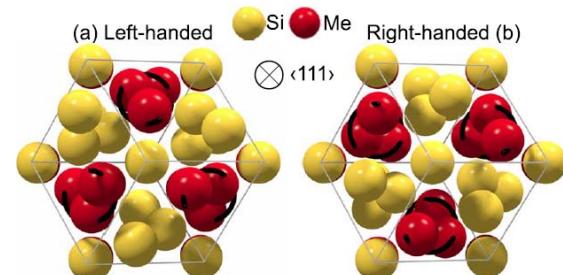
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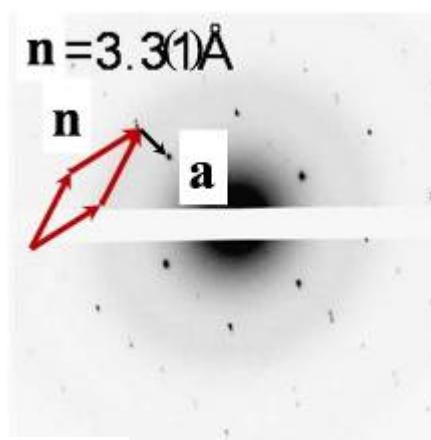
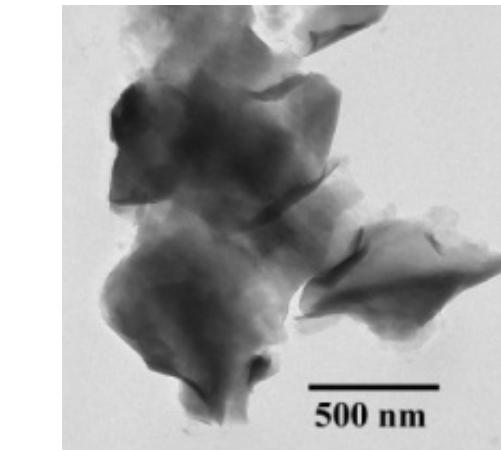
Quantum critical phenomena in the nano-world.
(Bad conductors, 2D systems).



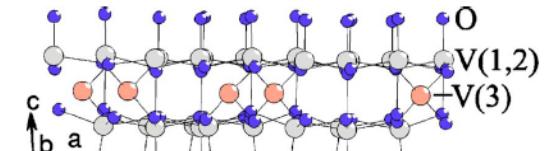
Quantum criticality in strongly correlated metals.
(Good conductors, 3D systems)



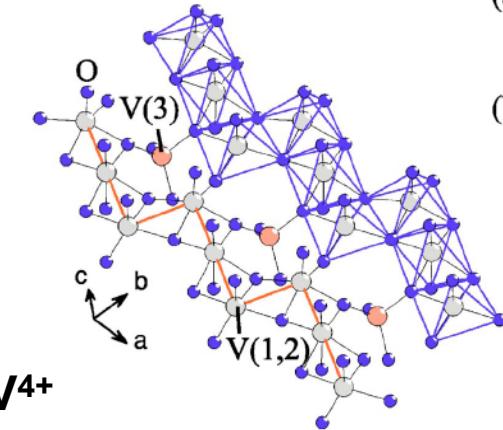
Final remarks



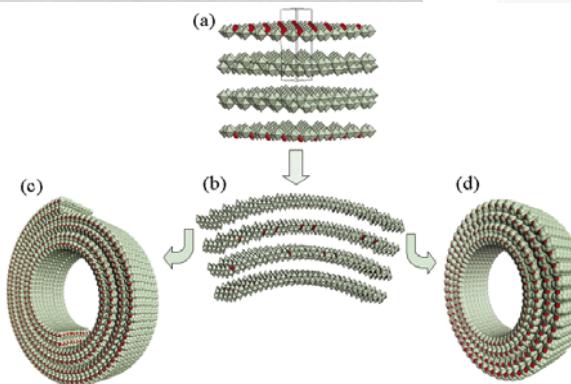
**VO_x nanolayers
(pre-VO_x-nanotubes)**



(a)

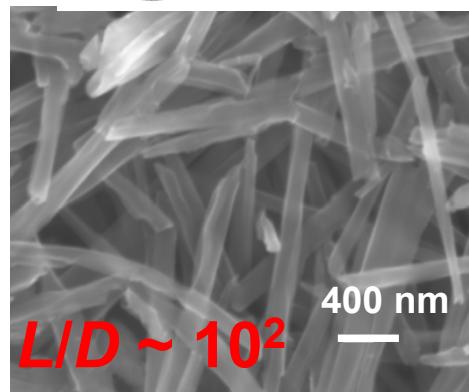


(b)

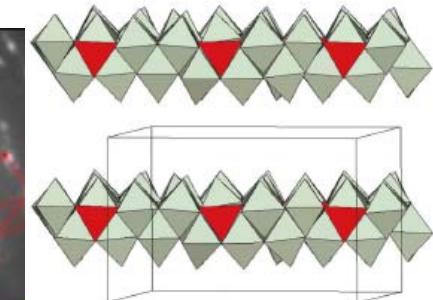
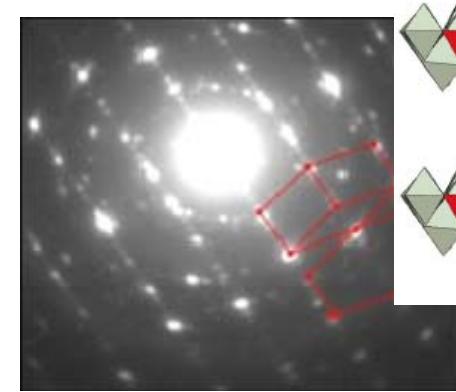
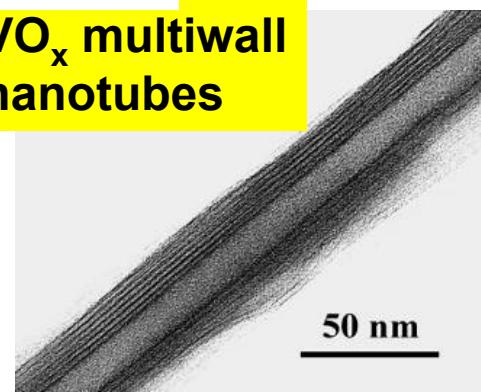


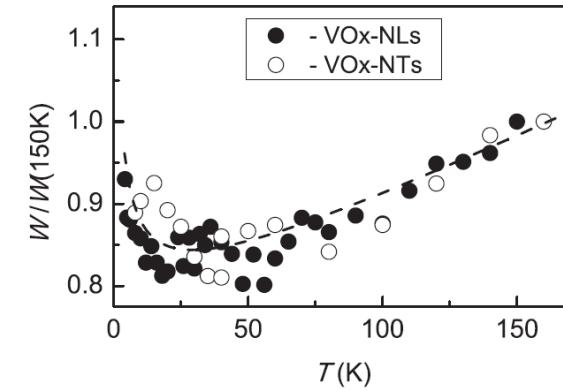
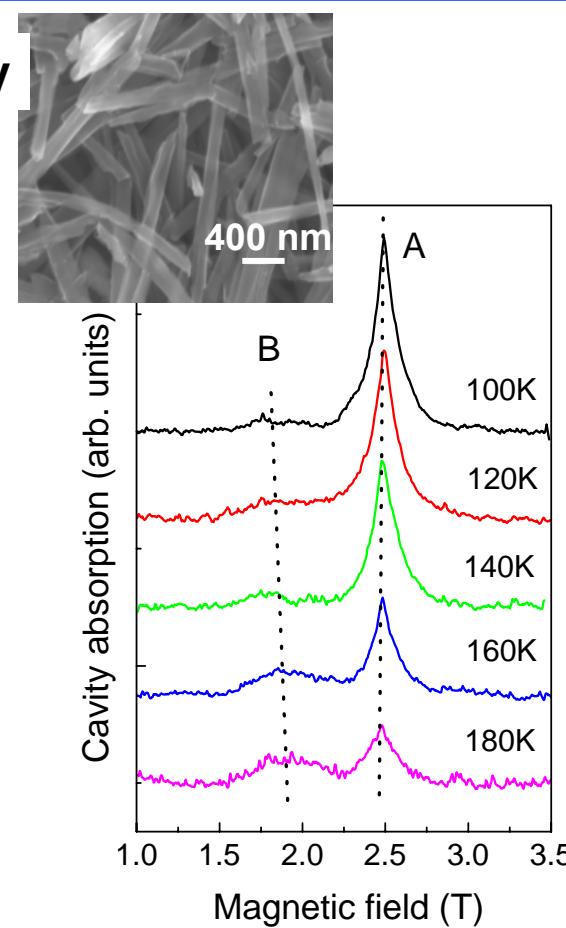
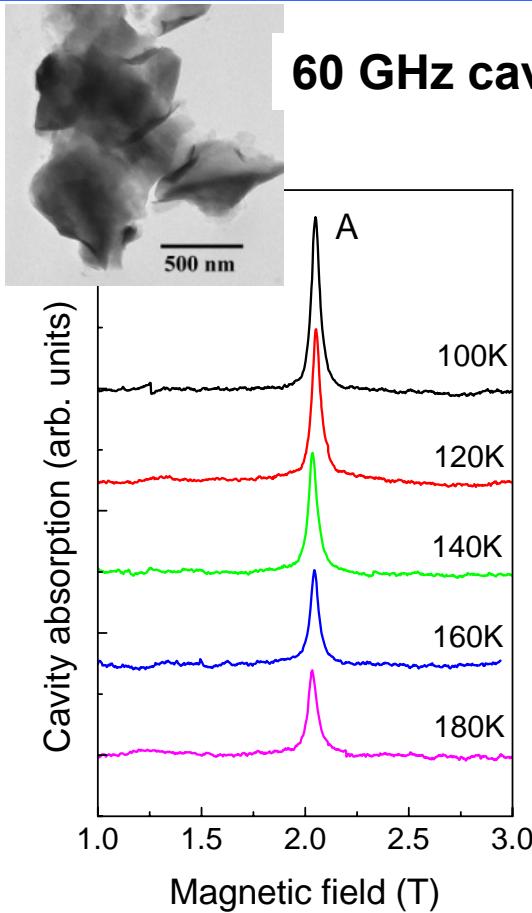
VO_x plane
6.13(1) Å

33 Å
V oxidation rate +4.2
80% V⁴⁺
20% V⁵⁺



**VO_x multiwall
nanotubes**

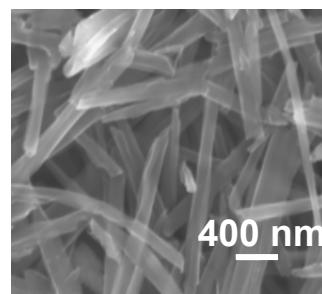
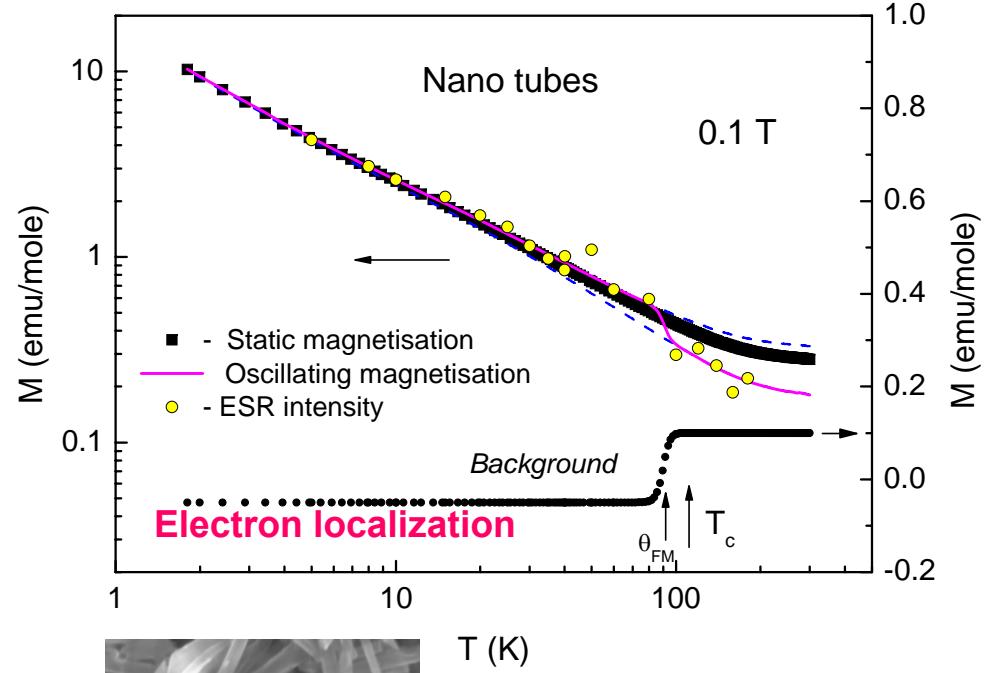
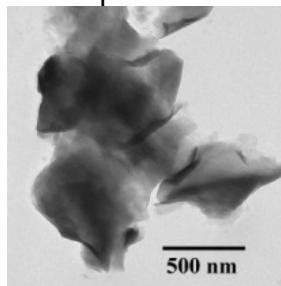
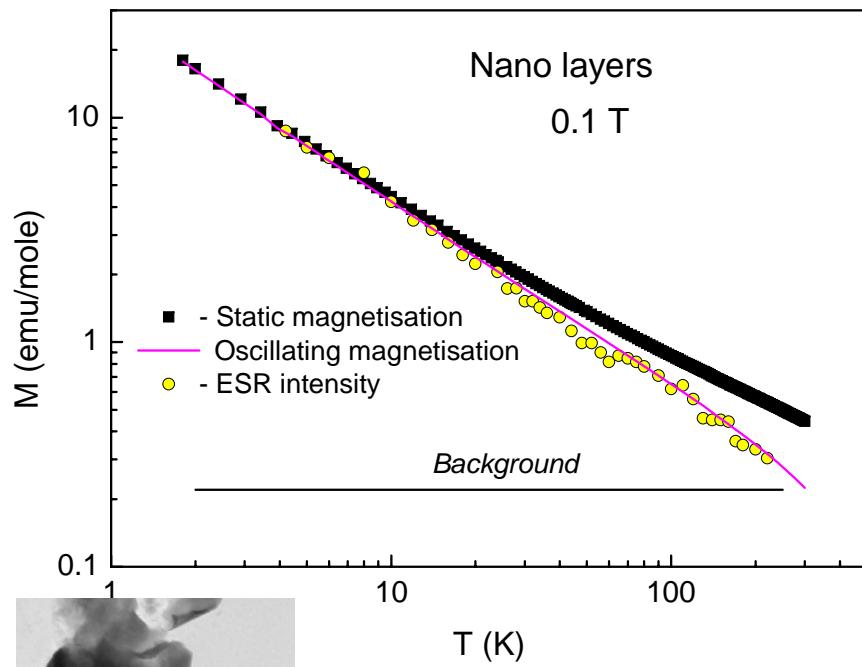




Again staggered field??
2D spin system rather
than 1D...

$g=1.96\pm0.02$ V^{5+} $S=0$ – non-magnetic ion
 V^{4+} $S=1/2$ – magnetic ion

$g_A=1.96\pm0.02$ V^{4+}
 $g_B=2.5\pm0.1$ $\text{V}^{4+}-\text{V}^{4+}$
AF dimer



$$M = M_{osc} + M_b$$

SQUID
(static)

ESR
(dynamic)

Background

Simple Van Vleck type
background structure.

$$\chi_b = \chi_{VanVleck} + \chi_{Pauli} + \chi_{Hubbard}$$



Change of background



free electrons in the sample

Absolute calibration

Curie constants for $V^{4+}(S=1/2)$ subsystem

Chemical analysis +
X-ray photoemission

X and Y numbers in $VO_x(C_{16}H_{33}NH_2)_Y$.
Average V charge ζ .

ESR

The model:

V^{4+} localized magnetic moments form dimers
and quasi free spins.
 g -factors for quasi free spins and dimers are
known.

$$\zeta = 5 - x = 5 - (x_l + x_e) = 5 - (x_f + x_d + x_e)$$

$$x_{empty} = 1 - (x_f + x_d + x_e)$$

$$C_f m = \frac{N_A \mu_B^2}{k_B} g_f^2 \frac{S(S+1)}{3} x_f = \frac{N_A \mu_B^2}{4k_B} g_f^2 x_f$$

$$x(V^{4+}) = x_f + x_d$$

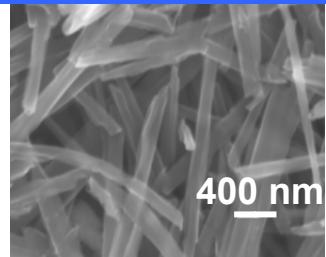
$$C_d m = \frac{N_A \mu_B^2}{k_B} g_d^2 x_d$$

x_f, x_d – electrons localized at V^{5+} making them V^{4+} , x_e - some other electrons,
different from those in V^{4+} state, m - molecular mass



Concentration estimates of various spin states.

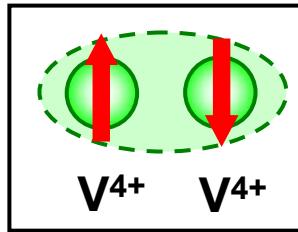
22



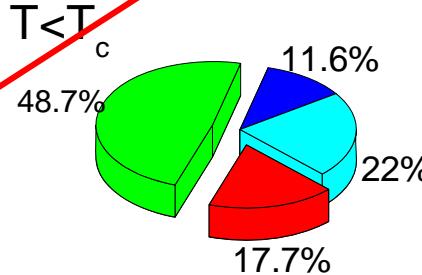
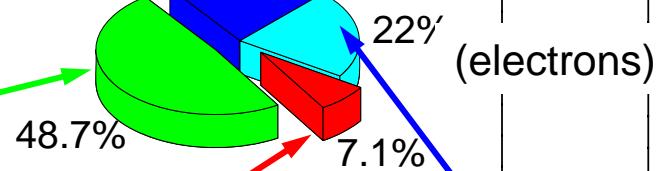
VO_x nanotubes

- 1- free spins
- 2- dimers
- 3- mixed state electrons
- 4- mixed state V⁵⁺

LMMs
state

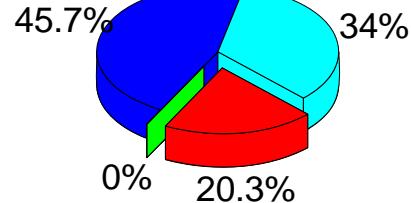
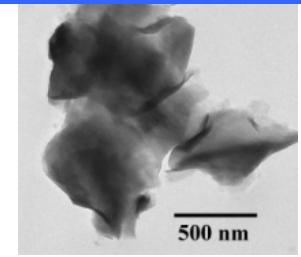


T>T_c 22.2% (empty V⁵⁺)

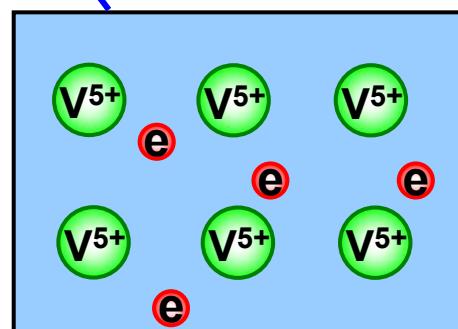


VO_x nano layers

- 1- free spins
- 2- dimers
- 3- mixed state electrons
- 4- mixed state V⁵⁺



“Mixed” state



Free electrons + V⁵⁺

or
electrons in upper
Hubbard band

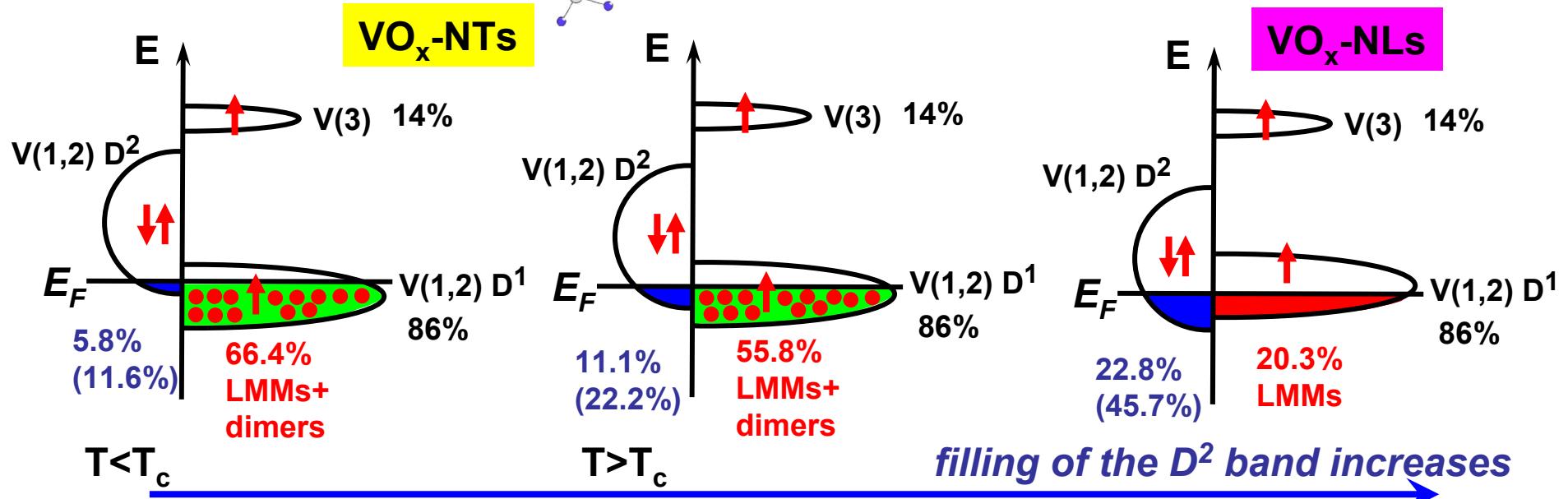
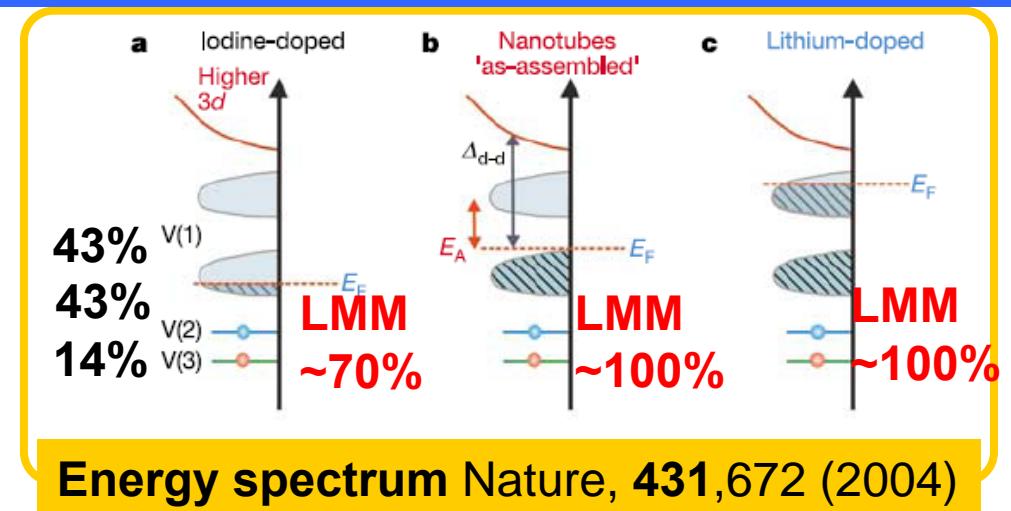
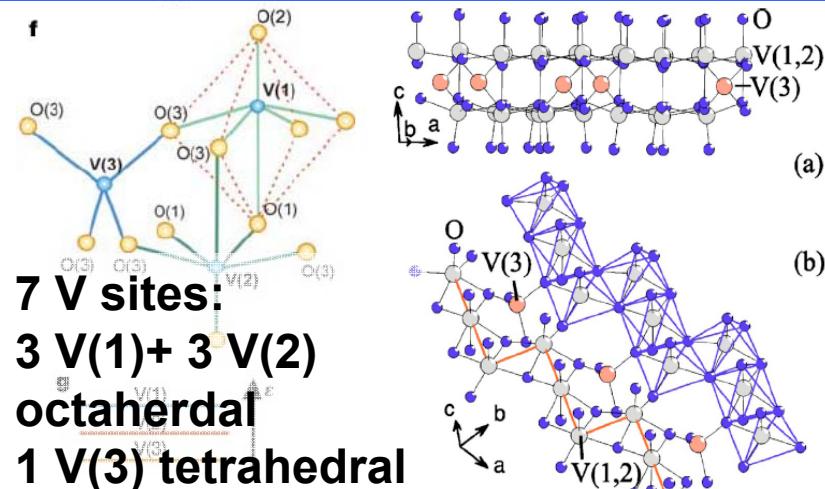
Scrolling of the VO_x layers:

- 1) Dimers;
- 2) Localization of electrons at low temperatures (T<T_c).



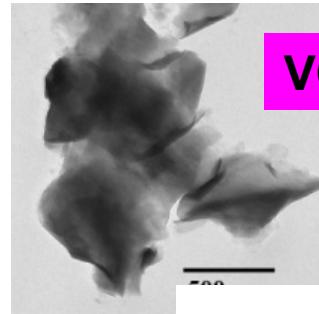
Concentration estimates and DOS of various spin states.

23

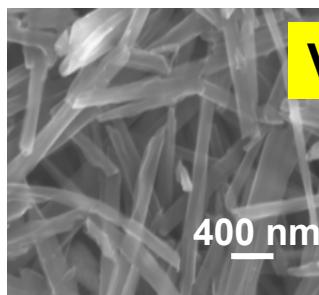
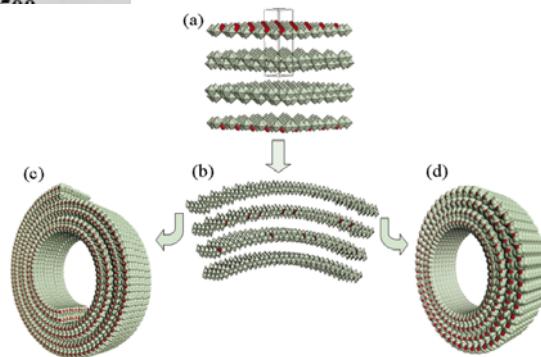


Magnetic transition= shift of the relative positions of D¹ and D².

Scrolling= enhancement of Hubbard repulsion

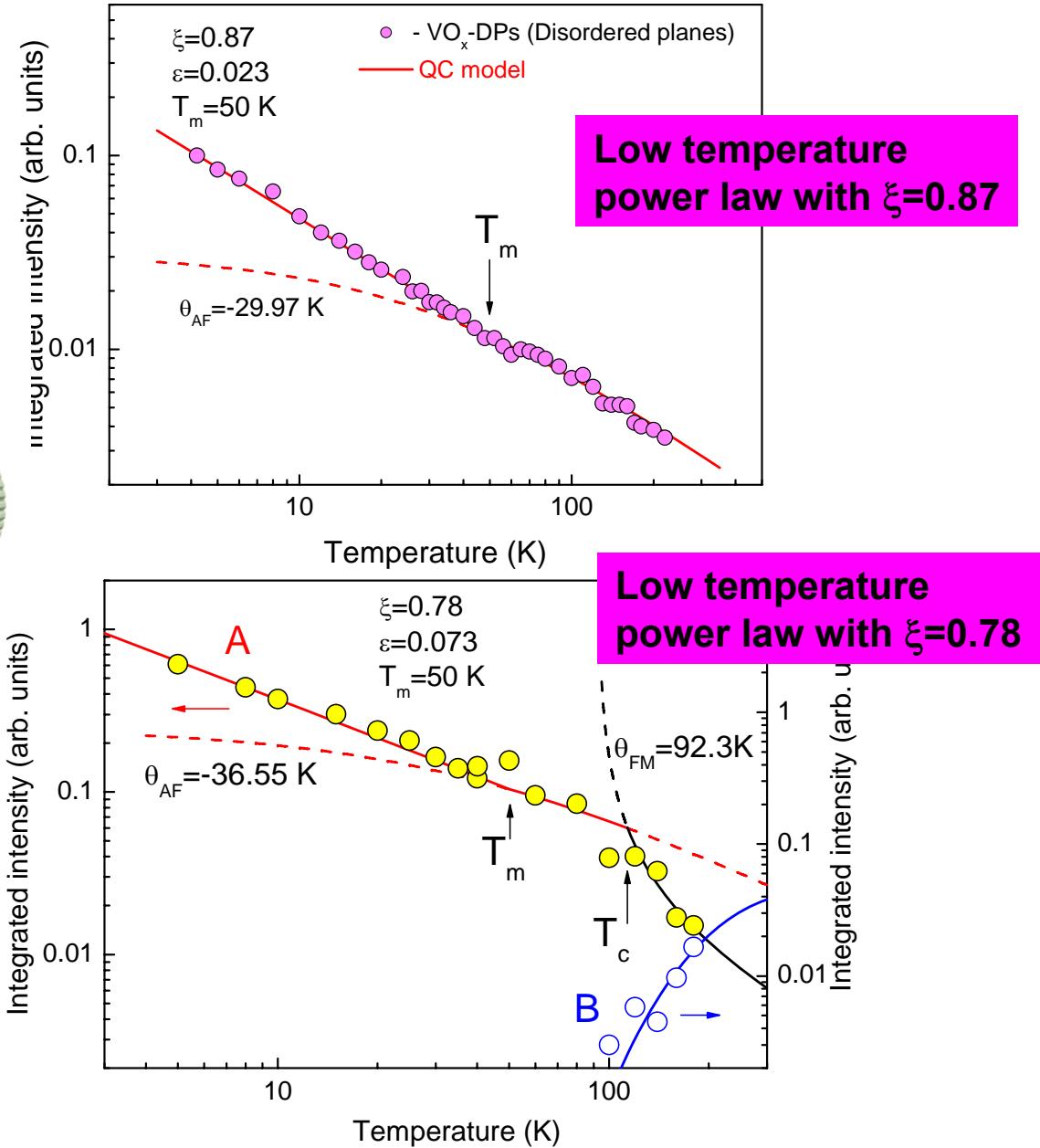


VO_x nanolayers

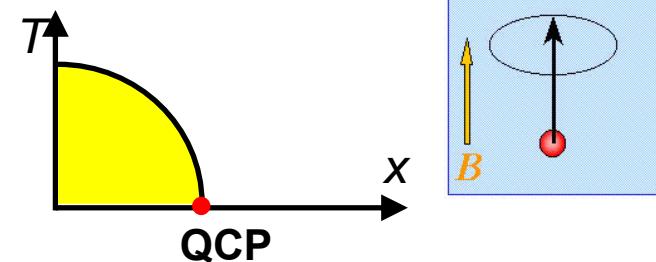


VO_x nanotubes

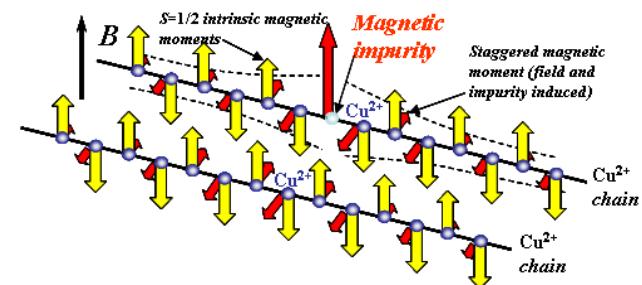
Good agreement between experiment and quantum critical model.



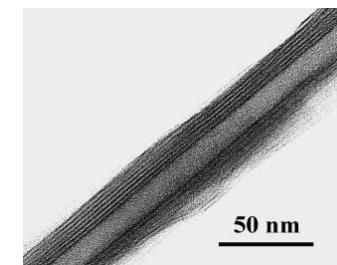
Introduction. Quantum criticality and electron spin resonance



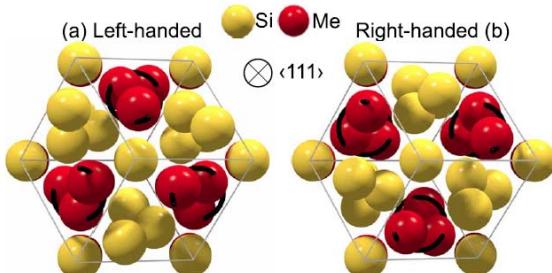
Quantum spin chains in disorder driven
quantum critical regime.
(Dielectrics, 1D systems)



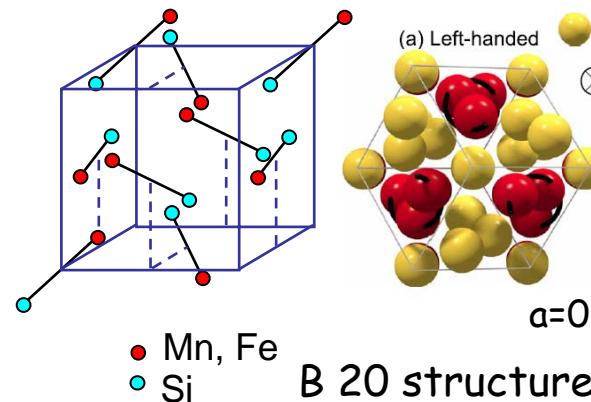
Quantum critical phenomena in the nano-world.
(Bad conductors, 2D systems)



Quantum criticality in strongly correlated metals.
(Good conductors, 3D systems)



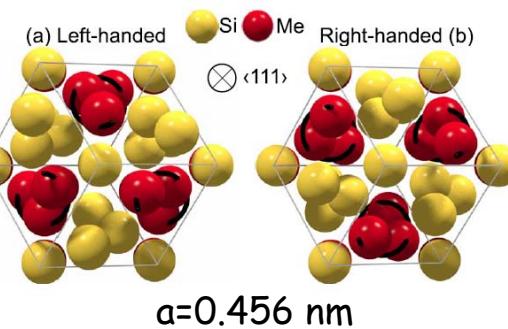
Final remarks



$Mn_{1-x}Fe_xSi$ substitutional solid solutions ($x < 0.3$).

After decay of the phase with long-range magnetic order an intermediate phase with short-range (fluctuation-driven) magnetic order should be formed. There is hidden QCP.

Intermediate phase may mask quantum critical anomalies like divergent magnetic susceptibility.



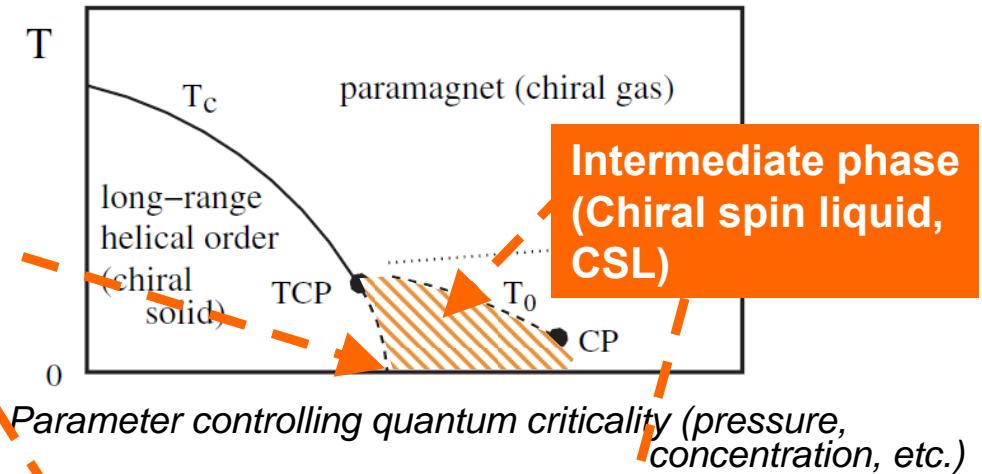
L 96, 047207 (2006)

PHYSICAL REVIEW LETTERS

week ending
3 FEBRUARY 2006

Blue Quantum Fog: Chiral Condensation in Quantum Helimagnets

Sumanta Tewari,^{1,2} D. Belitz,^{1,3} and T. R. Kirkpatrick^{1,2}



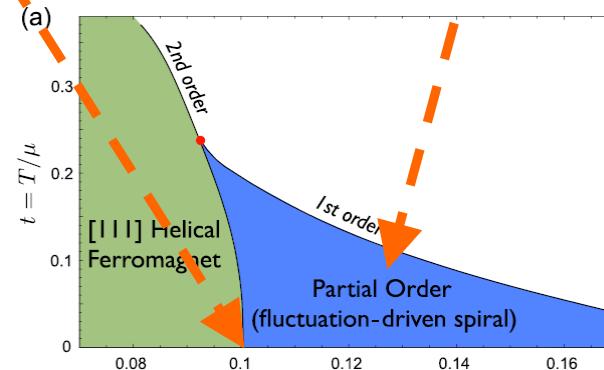
PRL 108, 067003 (2012)

PHYSICAL REVIEW LETTERS

week ending
10 FEBRUARY 2012

Quantum Order-by-Disorder Near Criticality and the Secret of Partial Order in MnSi

Frank Krüger,¹ Una Karahasanovic,¹ and Andrew G. Green²



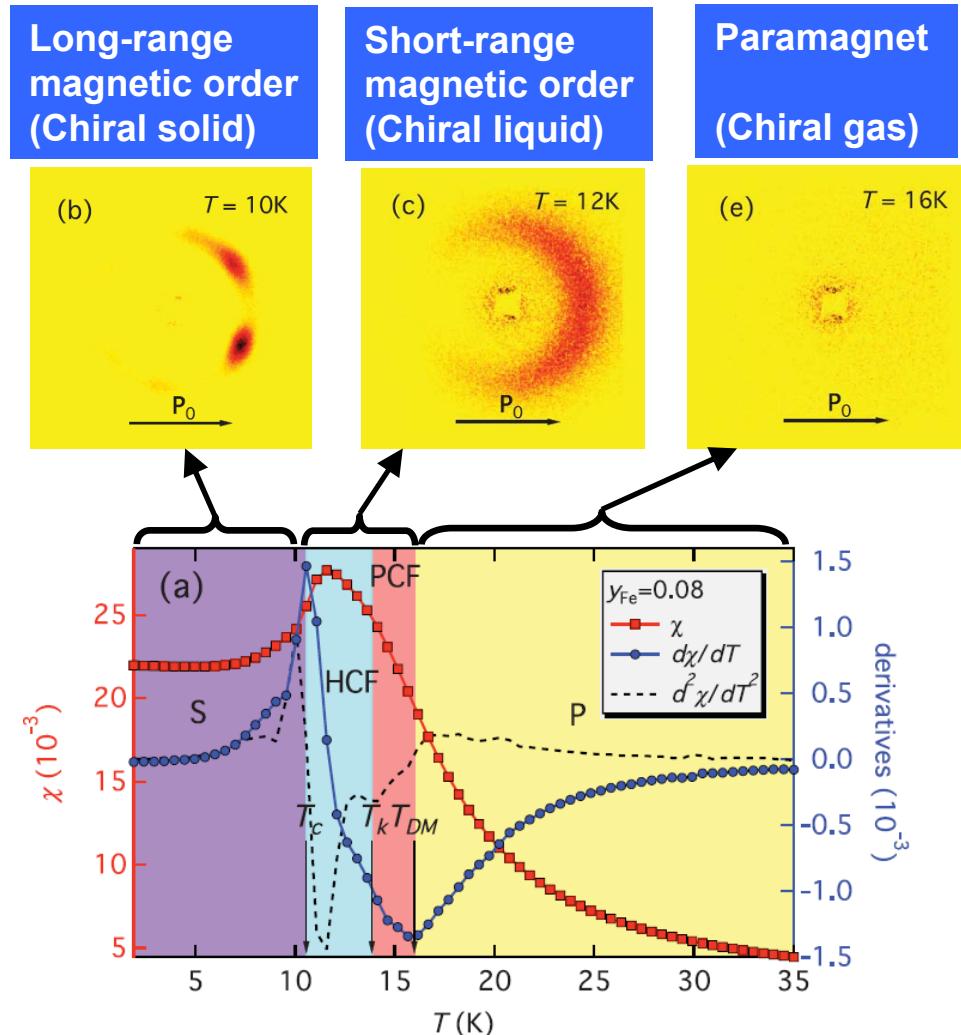
Parameter controlling quantum criticality (pressure, concentration, etc.)



PHYSICAL REVIEW B 83, 224411 (2011)

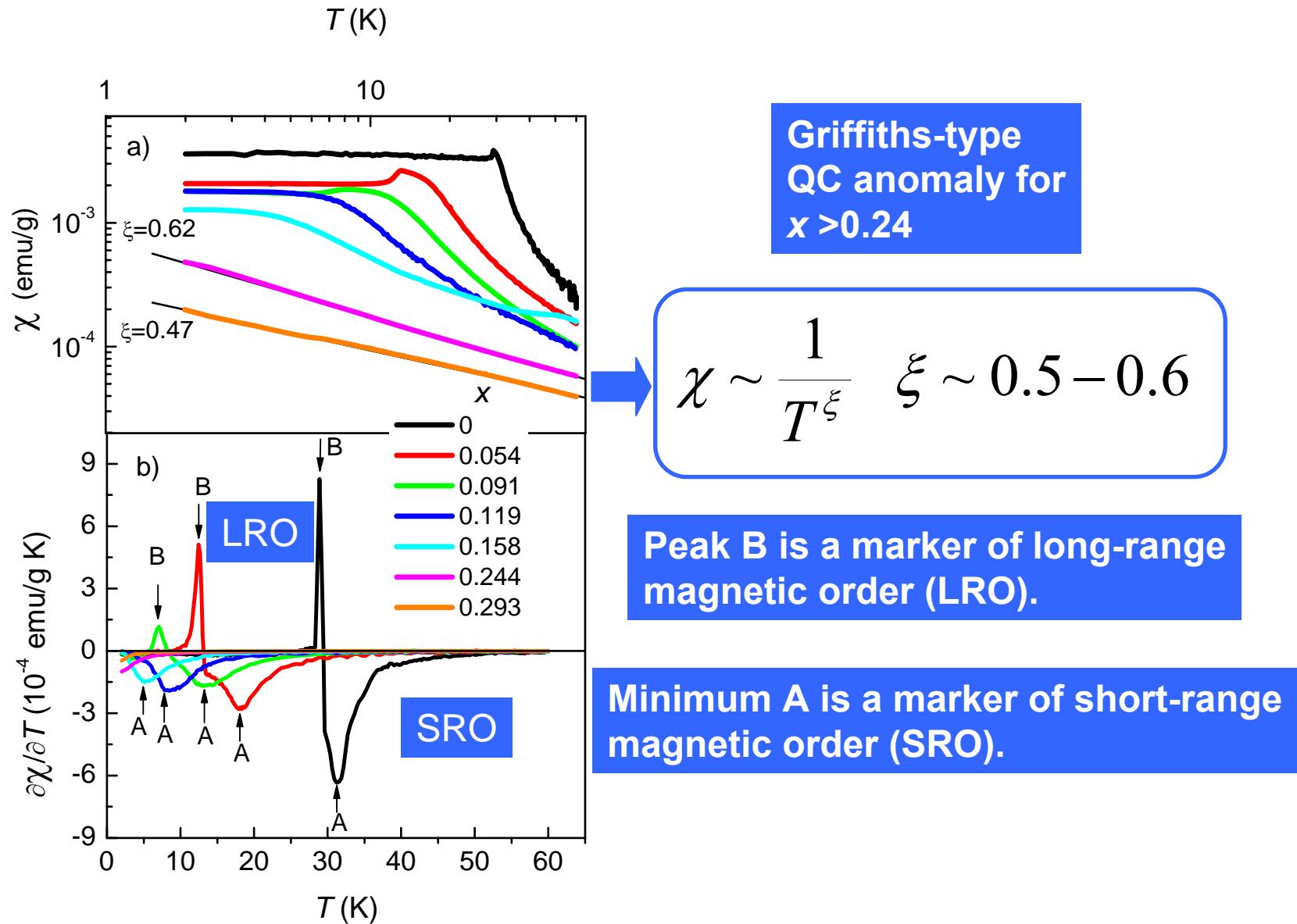
Chiral criticality in the doped helimagnets $\text{Mn}_{1-y}\text{Fe}_y\text{Si}$

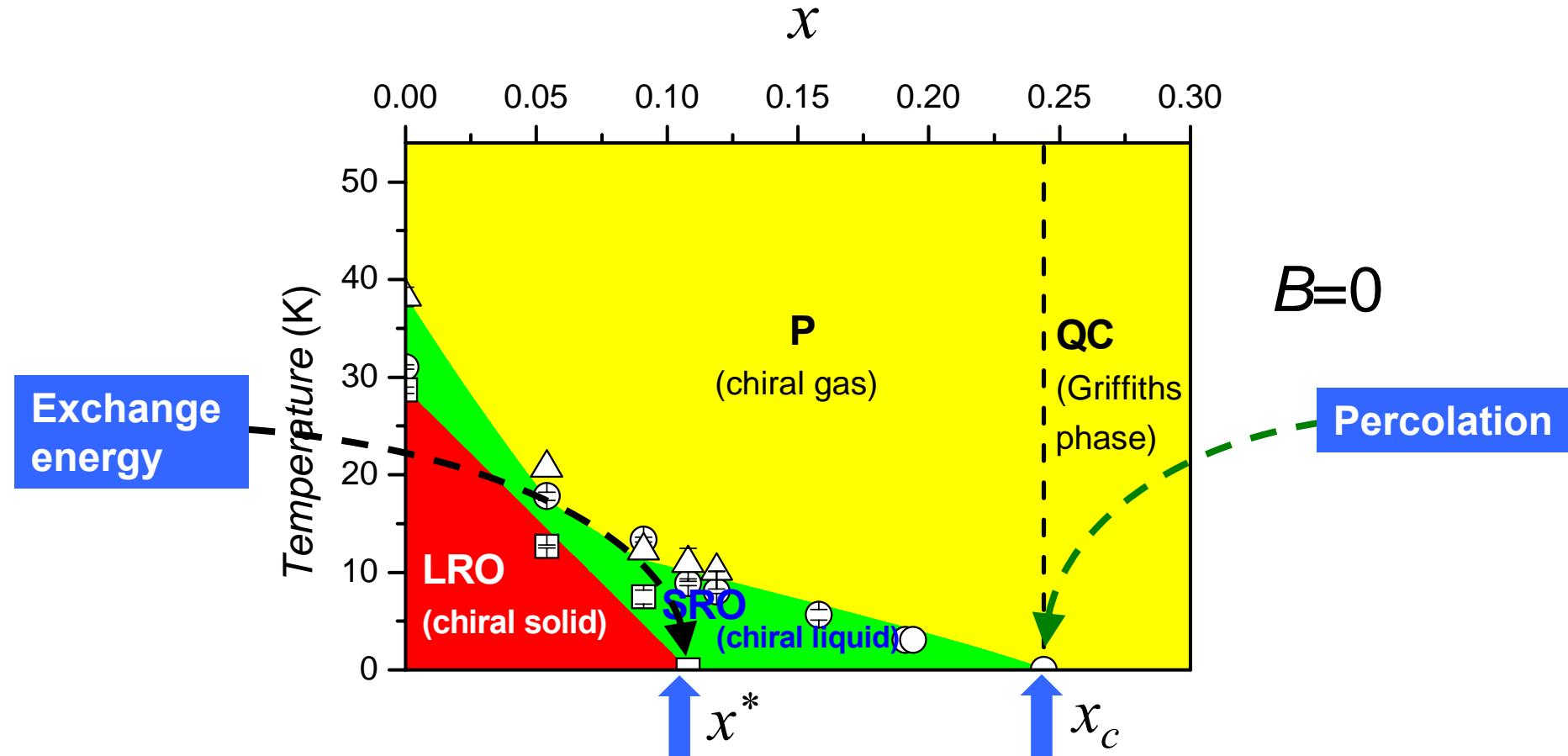
Sergey V. Grigoriev,¹ Evgeny V. Moskvin,¹ Vadim A. Dyadkin,¹ Daniel Lamago,^{2,3} Thomas Wolf,³ Helmut Eckerlebe,⁴ and Sergey V. Maleyev¹



Correlation between magnetic susceptibility and polarized neutron scattering data is established in $\text{Mn}_{1-x}\text{Fe}_x\text{Si}$.

The extrema of the $\partial\chi/\partial T$ derivative may be used for identification of the magnetic phases with long-range and short-range magnetic order.

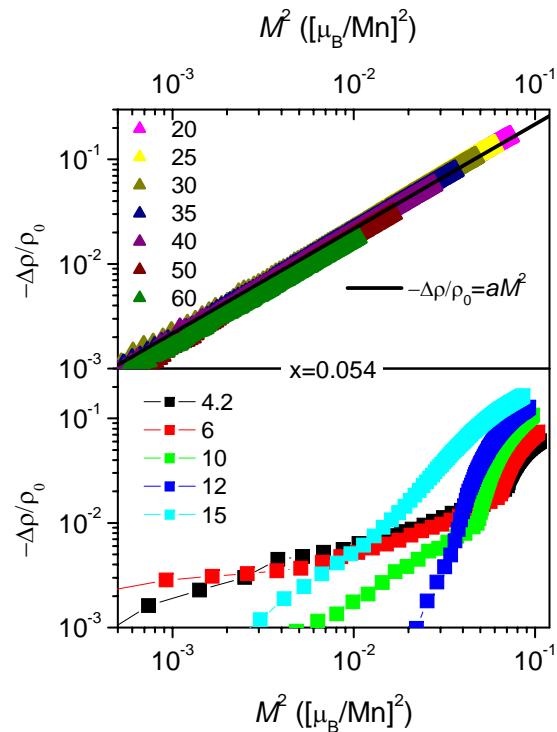




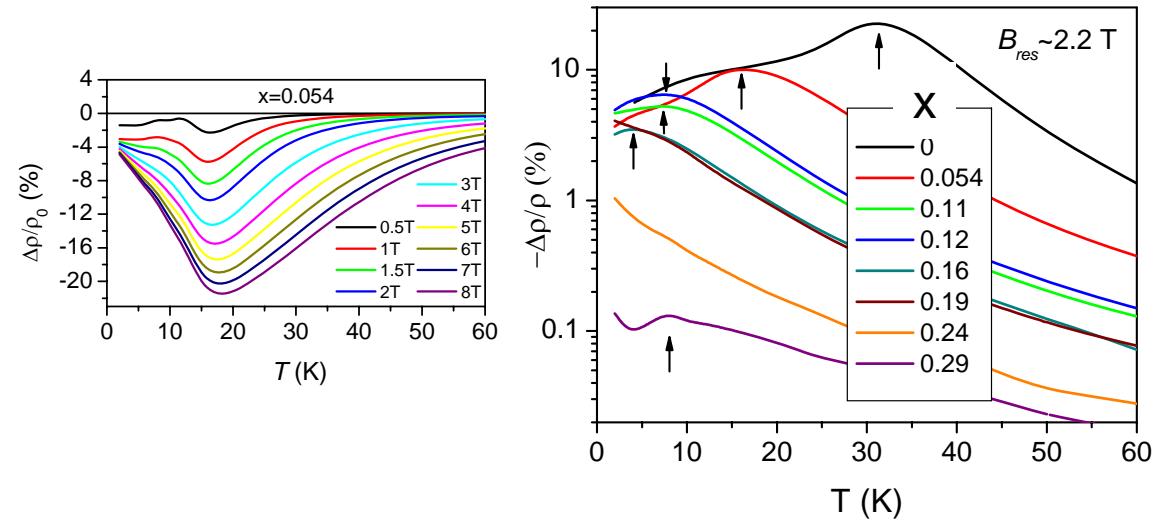
There are two quantum critical points, x^* and x_c . The first QC point $x^* \sim 0.11$ corresponds to disappearance of LRO and is a hidden one, which is located inside the SRO phase. The second QC point $x_c \sim 0.24$ is a “true” one and marks suppression of the magnetic phase with SRO (chiral spin liquid).



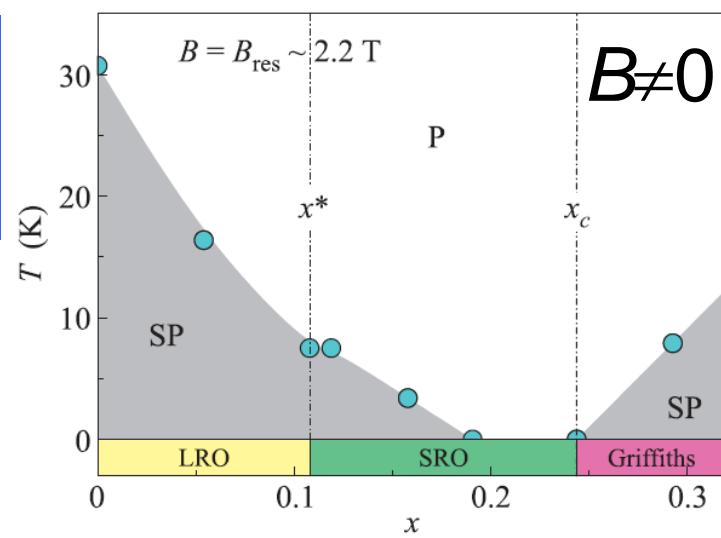
Yosida mechanism holds
in $Mn_{1-x}Fe_xSi$.



Magnetoresistance maximum evolution
with iron concentration.



Phase diagram
in magnetic field
corresponding to
magnetic resonance



Pis'ma v ZhETF, vol. 100, iss. 1, pp. 30–33

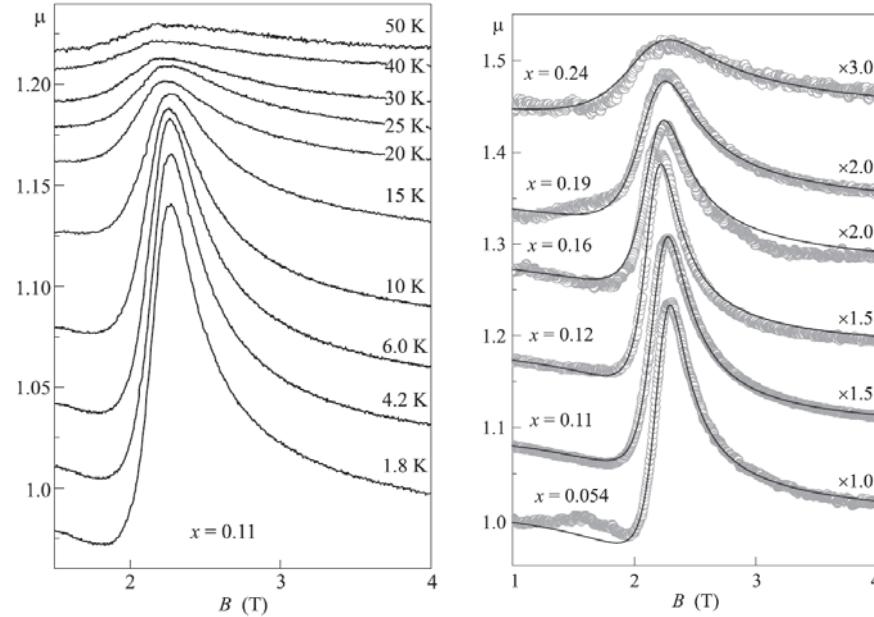
© 2014 July 10

Anomalous spin relaxation and quantum criticality
in $Mn_{1-x}Fe_xSi$ solid solutions¹⁾

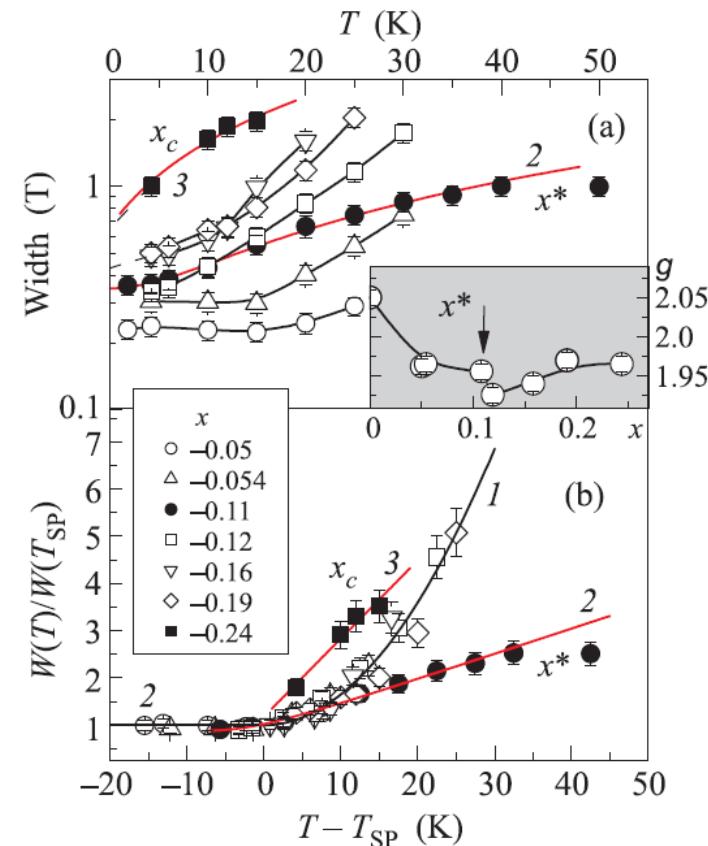
S. V. Demishev⁺⁺²⁾, A. N. Samarin^{++*}, V. V. Glushkov^{++*}, M. I. Gilmanov^{++*}, I. I. Lobanova^{++*}, N. A. Samarin⁺,
A. V. Semeno⁺, N. E. Sluchanko⁺, N. M. Chubova[×], V. A. Dyadkin[×], S. V. Grigoriev[×]



60 GHz ESR spectra



**Strong broadening of the line width
(enhancement of spin fluctuations)
with iron concentration.**



Universal scaling $W(T)/W(T_{SP})=1+a(T-T_{SP})^2$ for all concentrations except quantum critical points x^* and x_c .

Violation of the standard Korringa relaxation law $W(T) \sim 1/\chi(T) \sim (T-T_{SP})$.

Weakening of the $W(T)$ temperature dependence just at quantum critical points.



Theory of the ESR in strongly correlated metals

Korringa relaxation should be valid even in strongly correlated metallic case

P. Schlottmann, Phys. Rev. B **86**, 075135 (2012)

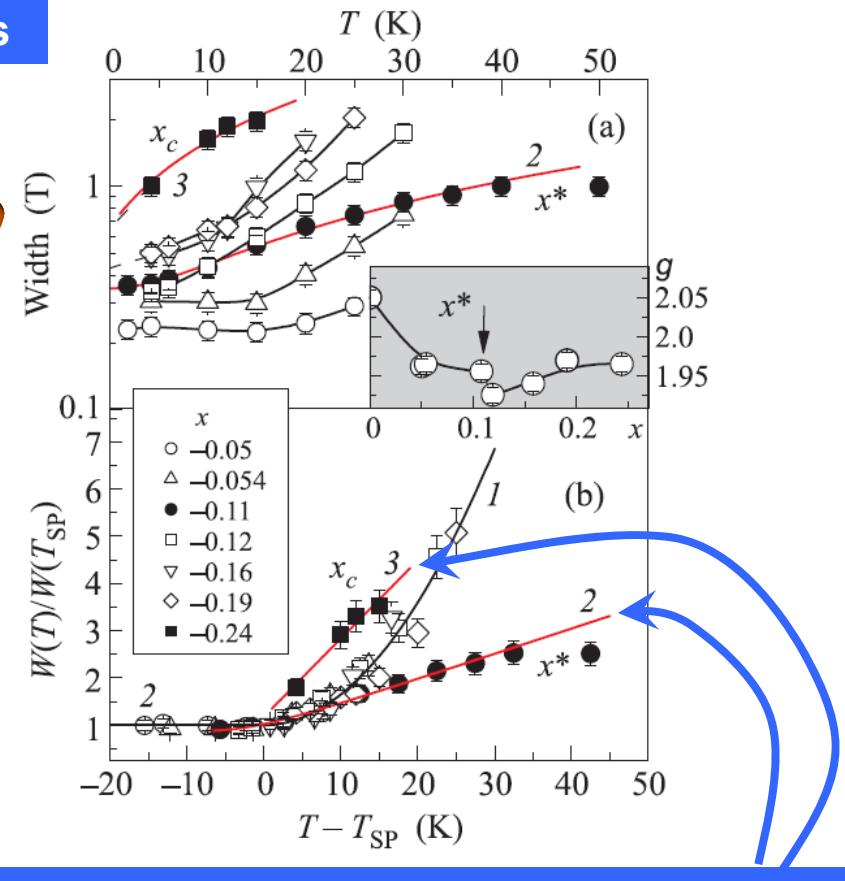


Expression for ESR line width in strongly correlated metal in vicinity of the QC point:

$$W(T) = AT \arctan(T/T_x) + W_0$$

T_x is a crossover temperature between Fermi-liquid ($T \ll T_x$) and non-Fermi-liquid ($T \gg T_x$) regimes.

P. Wölfe and E. Abrahams, Phys. Rev. B **80**, 235112 (2009).



Reasonable fit of the experimental data assuming $T_x \sim T_{SP}$, i.e. $T_x \sim 11$ K for x^* and $T_x \sim 0$ K for x_c .

Even in strong magnetic field QC points x^* and x_c derived in the limit $B \rightarrow 0$ are still here.

ESR is a right tool to visualize QC points including hidden QC point.

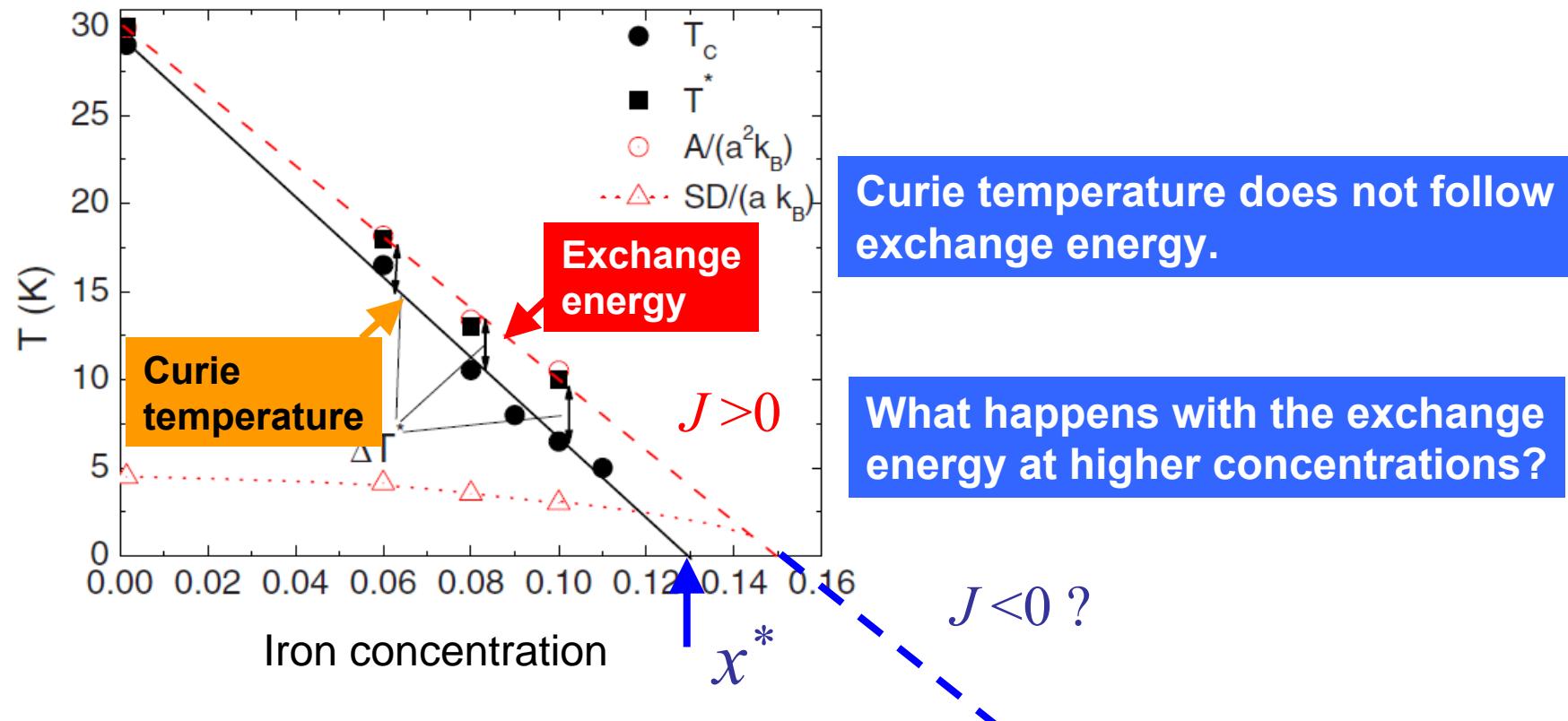




PHYSICAL REVIEW B 79, 144417 (2009)

Helical spin structure of $Mn_{1-y}Fe_ySi$ under a magnetic field: Small angle neutron diffraction study

S. V. Grigoriev,¹ V. A. Dyadkin,¹ E. V. Moskvin,^{1,2} D. Lamago,^{3,4} Th. Wolf,⁴ H. Eckerlebe,⁵ and S. V. Maleyev¹



Curie temperature does not follow exchange energy.

What happens with the exchange energy at higher concentrations?

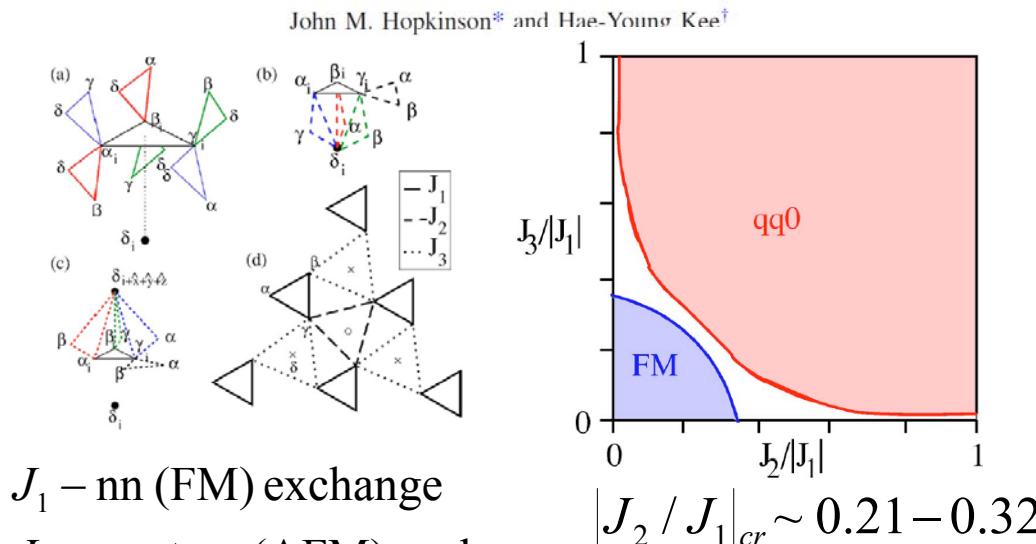
Why Mn subsystem breaks into spin clusters at relatively low concentrations, where the formal percolation may exist ($x_c \sim 0.83$ instead of observed $x_c \sim 0.24$) ?



Why intermediate phase is formed? Why LRO phase is suppressed?

PHYSICAL REVIEW B 75, 064430 (2007)

Microscopic model for spiral ordering along (110) on the MnSi lattice

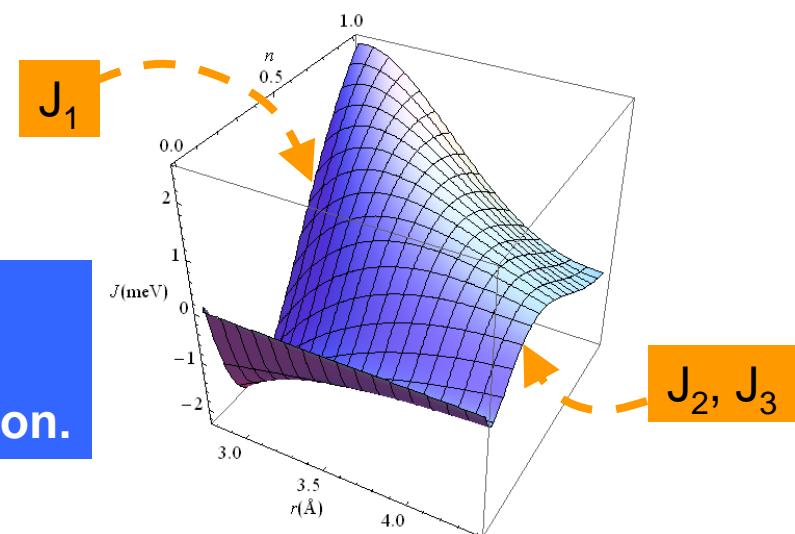


In Heisenberg paradigm RKKY exchange defines J_1, J_2, J_3 parameters, which may be tuned by variation of the electron concentration.

Frustration try to align spirals along (110).

DM interaction try to align spirals along (111)

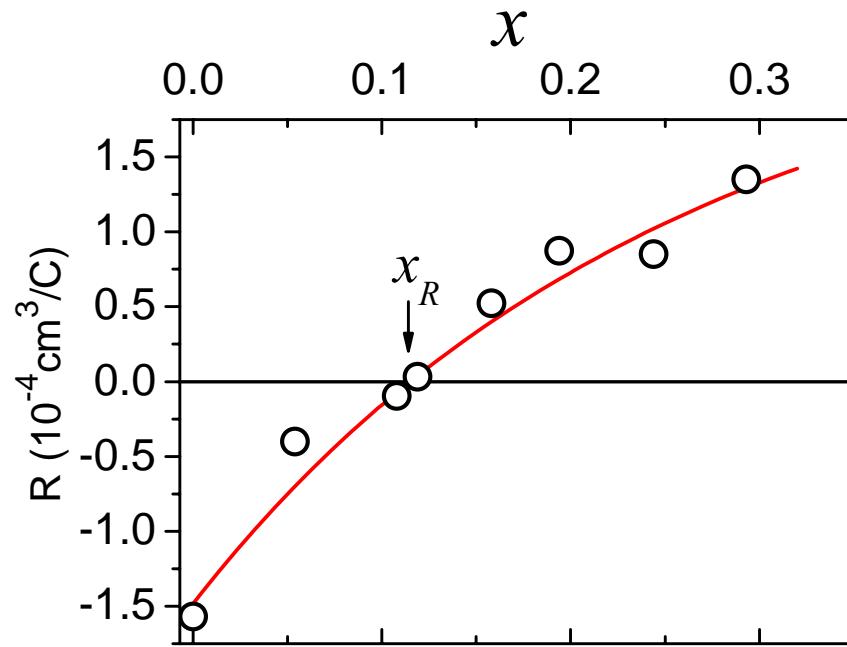
Competition between two interactions may lead to loosing of the long-range order and formation of the chiral liquid state.



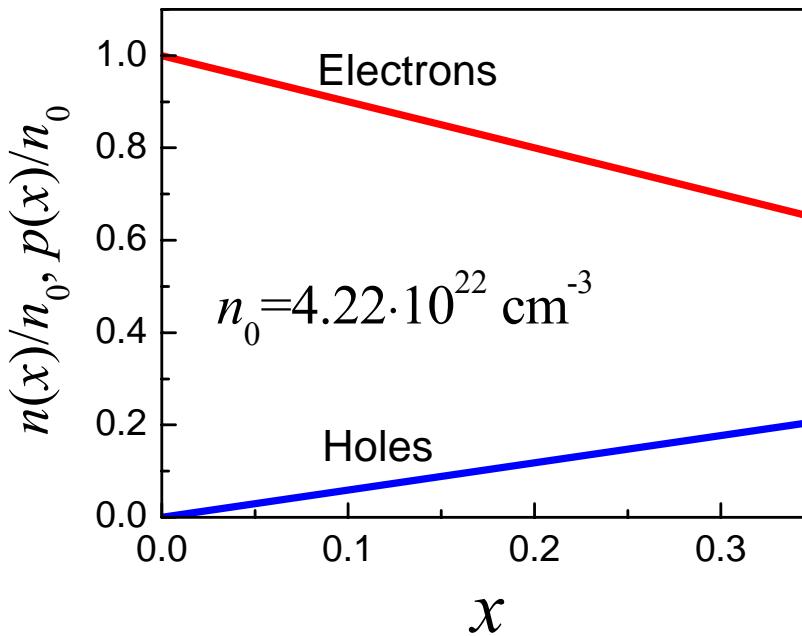


Ordinary Hall effect coefficient in $Mn_{1-x}Fe_xSi$

(Glushkov, Lobanova, 2015)



The model:



Two groups of charge carriers: electrons and holes

$$x_R = 0.115 \sim x^*$$

$$|\mu_h / \mu_e| = 0.28$$

Electrons are from Mn, holes are from Fe.



$$J_{RKKY}(r) = \text{const} \frac{m}{r^4} \varphi(2k_F r)$$

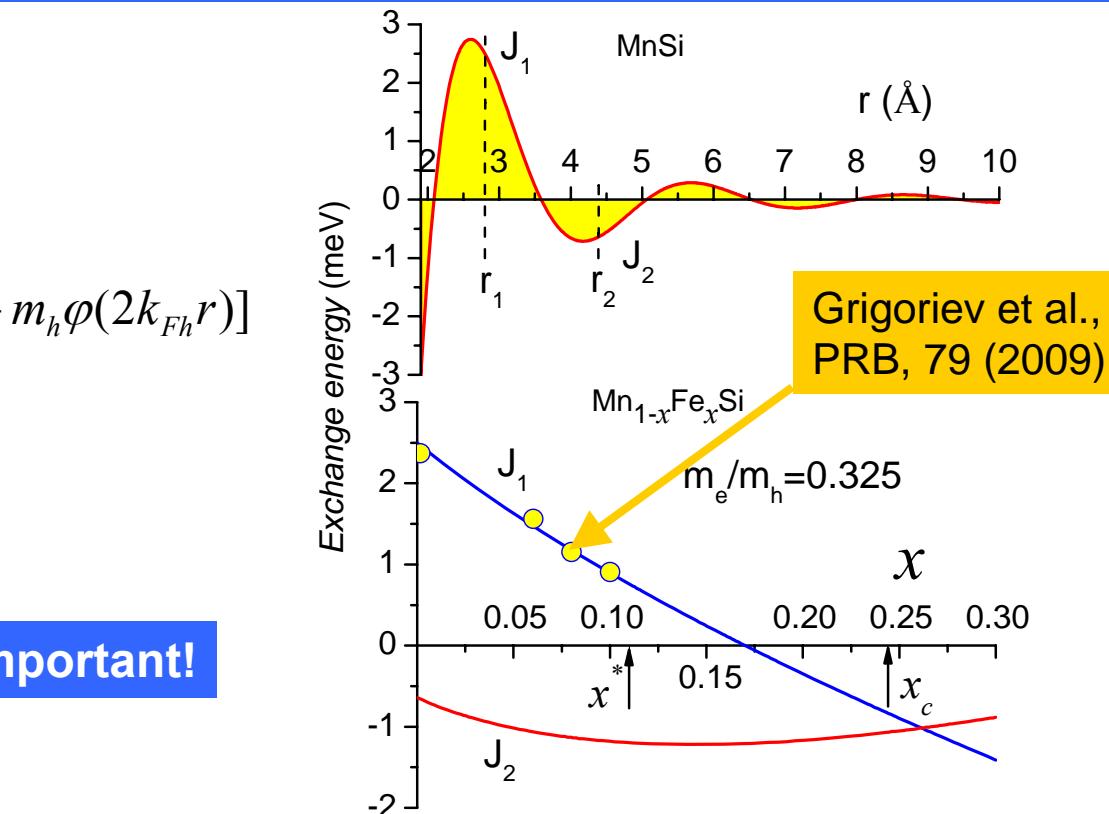
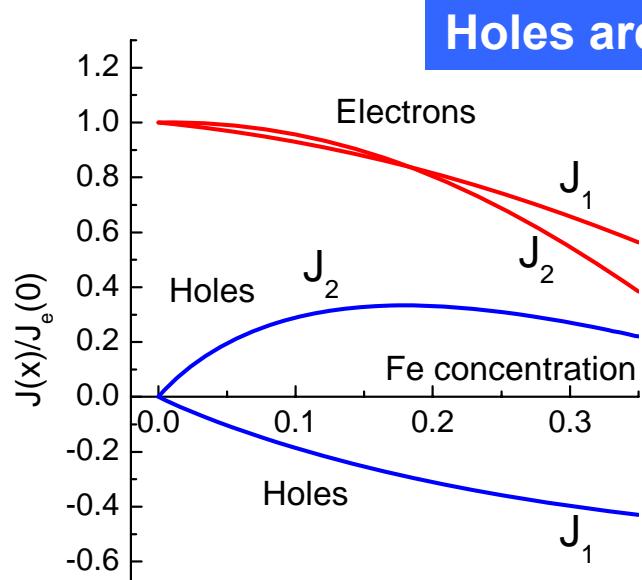
$$\varphi(x) = x \cos(x) - \sin(x)$$

$$J = J_e + J_h = \frac{\text{const}}{r^4} [m_e \varphi(2k_{Fe} r) + m_h \varphi(2k_{Fh} r)]$$

$$k_{Fe}(x) = k_{Fe}(0) [n(x)/n_0]^{1/3}$$

$$k_{Fh}(x) = k_{Fe}(0) [p(x)/n_0]^{1/3}$$

$$k_{Fe}(0) = (3\pi^2 n_0)^{1/3}$$



In the vicinity of QC points $|J_1| \sim |J_2|$.
Frustration is essential.

Frustration: T_c goes to zero faster than J_1 .

Frustration: breaks coherence of the magnetic state of Mn ions and “helps” forming separated spin clusters at x_c .



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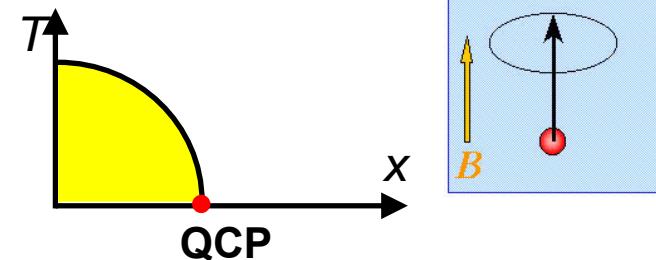
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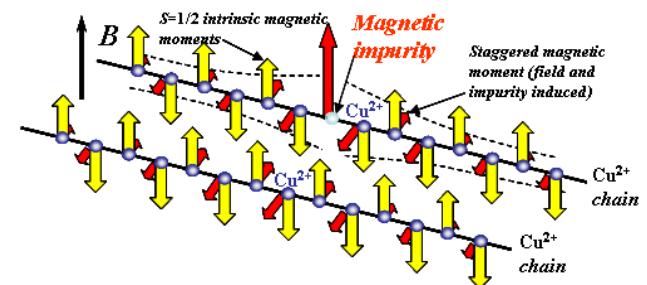
"Scientists now believe that there may be a connection between the recent earthquakes and the record snowfall."

Quantum criticality in $Mn_{1-x}Fe_xSi$ is driven by change of electrons and holes concentration (i.e. change of the Fermi surface) together with disorder effects.

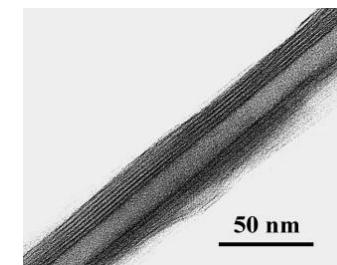
Introduction. Quantum criticality and electron spin resonance



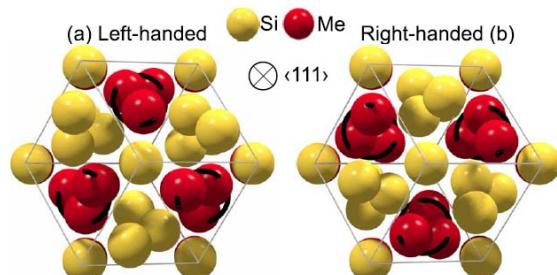
Quantum spin chains in disorder driven,
quantum critical preline.
(Dielectrics, 1D systems).



Quantum critical phenomena in the nano-world.
(Bad conductors, 2D systems)



Quantum criticality in strongly correlated metals.
(Good conductors, 3D systems)



Final remarks



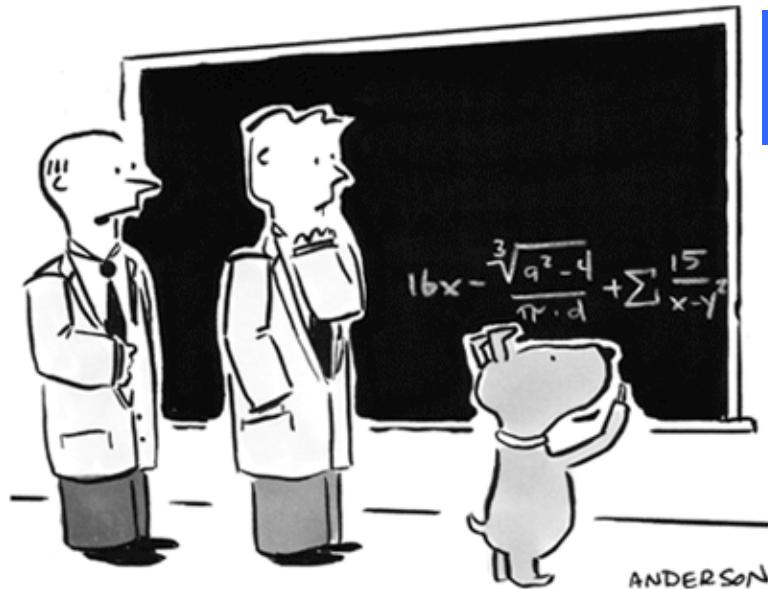


ESR is a right tool to study quantum criticality. Are other tools right?

**Look for staggered field (low temperature growth of the line width)
and you will find QC phenomena!**

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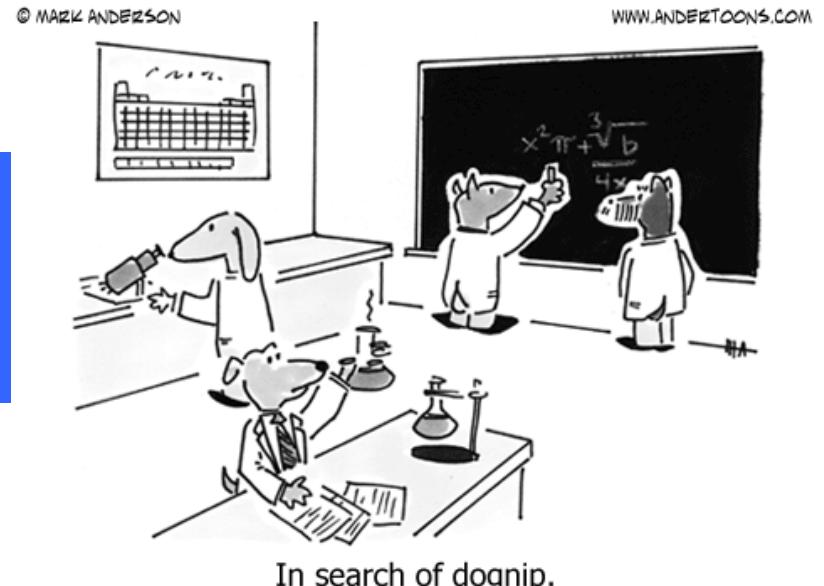
"I think he's trying to tell us something."

**ESR in metals is almost unexplored route
to quantum criticality.**

**In many cases quantum criticality is
nothing but a section of nanomagnetism.**

Special thanks to my colleagues and co-authors:
A.V.Semeno, N.E.Sluchanko, V.V.Glushkov, T.V.Ischenko, N.A.Samarin,
I.I.Lobanova, A.N.Samarin, A.L.Chernobrovkin (GPI),
A.V.Grigorieva, E.A.Goodilin (MSU),
S.V.Grigoriev, N.Chubova (PNPI), H.Ohta (Kobe University)

**Thank you for your
attention!**



Acknowledgements

Programmes of Russian Academy of Sciences
"Electron spin resonance, spin-dependent electronic effects and spin technologies",
"Electron correlations in strongly interacting systems" and RFBR grant 13-02-00160

Theory

ISSN 1063-7834, Physics of the Solid State, 2009, Vol. 51, No. 3, pp. 547–551. © Pleiades Publishing, Ltd., 2009.
Original Russian Text © S.V. Demishev, 2009, published in Fizika Tverdogo Tela, 2009, Vol. 51, No. 3, pp. 514–517.

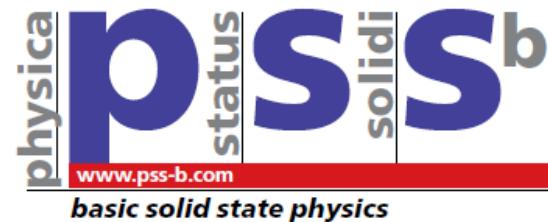
MAGNETISM
AND FERROELECTRICITY

Physics of the Solid State,
51 (2009)

Modeling of Magnetic Susceptibility of an Antiferromagnetic System with Disorder-Driven Quantum Critical Behavior

S. V. Demishev

Phys. Status Solidi B 247, No. 3, 676–678 (2010) / DOI 10.1002/pssb.200983003



Magnetic susceptibility of an antiferromagnetic system with disorder-driven quantum critical behavior

Sergey Demishev*

PHYSICAL REVIEW B 84, 094426 (2011)

PRB, 84 (2011)

Magnetic properties of vanadium oxide nanotubes and nanolayers

S. V. Demishev,^{1,*} A. L. Chernobrovkin,^{1,2} V. V. Glushkov,¹ A. V. Grigorieva,³ E. A. Goodilin,³ H. Ohta,⁴ S. Okubo,⁴ M. Fujisawa,⁴ T. Sakurai,⁵ N. E. Sluchanko,¹ N. A. Samarin,¹ and A. V. Semeno¹

CuGeO₃:Fe

JETP Letters, Vol. 73, No. 1, 2001, pp. 31–34. Translated from *Pis'ma v Zhurnal Èksperimental'noi Teoreticheskoi Fiziki*, Vol. 73, No. 1, 2001, pp. 36–40.
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JETP Jett., 73 (2001)

New Scenario for the Decay of Spin-Peierls State in CuGeO₃ : Fe. Onset of a Quantum Critical Point

S. V. Demishev*, R. V. Bunting, L. I. Leonyuk†, E. D. Obraztsova, A. A. Pronin,
N. E. Sluchanko, N. A. Samarin, and S. V. Terekhov

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Quantum Criticality and Collective Effects in Low-Dimensional Magnet CuGeO₃:Fe Probed by High Frequency EPR

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Oshikawa-Affleck theory

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**Anomalous temperature dependence of the ESR linewidth
in CuGeO₃ doped with magnetic impurities
and the universal relations in the Oshikawa-Affleck theory**

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**The Competition between Staggered Field and Antiferromagnetic
Interactions in CuGeO₃:Fe**

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CuGeO₃:Co

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Microwave EPR Spectroscopy of Cobalt-Doped Germanium Cuprate

S. V. Demishev^{1,4}, A. V. Semeno¹, N. E. Sluchanko¹, N. A. Samarin¹, A. A. Pronin¹,
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JETP Jett., 84 (2006)

New Polarization Effect and Collective Excitation in $S = 1/2$ Quasi-one-dimensional Antiferromagnetic Quantum Spin Chain[¶]

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CuGeO₃:Mn



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Quantum critical behavior induced by Mn impurity in CuGeO₃

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ESR Probing of Quantum Critical Phenomena in Doped $S = 1/2$ AF Quantum Spin Chain*

S. V. Demishev¹, A. V. Semeno¹, N. E. Sluchanko¹, N. A. Samarin¹,
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VO_x nanomaterials



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Electron spin resonance
and quantum critical phenomena
in VO_x multiwall nanotubes



phys. stat. sol. (RRL), 1-3 (2008)

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FM–AFM Crossover in Vanadium Oxide Nanomaterials

S. V. Demishev^{a,*}, A. L. Chernobrovkin^a, V. V. Glushkov^a, A. V. Grigorieva^b,
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PHYSICAL REVIEW B 84, 094426 (2011)

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Magnetic properties of vanadium oxide nanotubes and nanolayers

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Mn_{1-x}Fe_xSi system

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Is MnSi an Itinerant-Electron Magnet? Results of ESR Experiments

**S. V. Demishev*, A. V. Semeno, A. V. Bogach, V. V. Glushkov, N. E. Sluchanko,
N. A. Samarin, and A. L. Chernobrovkin**

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Magnetic phase diagram of MnSi in the high-field region

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N. A. Samarin, N. E. Sluchanko, V. M. Zimin, and A. V. Semeno**

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Quantum Bicriticality in Mn_{1-x}Fe_xSi Solid Solutions: Exchange and Percolation Effects[¶]

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Anomalous Spin Relaxation and Quantum Criticality in Mn_{1-x}Fe_xSi Solid Solutions[¶]

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